

EINSTEIN EQUATIONS FOR A GENERALIZED LAGRANGE SPACE OF ORDER 2 IN INVARIANT FRAMES

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Abstract

The study of higher order Lagrange spaces founded on the notion of bundle of velocities of order k has been given by Radu Miron and Gheorghe Atanasiu in [2].

The bundle of accelerations correspond in this study to $k=2$.

The notion of invariant geometry of order 2 was introduced by the author in [4].

In this paper we shall give the Maxwell equations of a generalized Lagrange space of order 2 in invariant frames.

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1. General Invariant Frames

Let us consider the bundle $E = Osc^2M$, a nonlinear connection N with the coefficients

$$\begin{pmatrix} N_{(1)j}^i & N_{(2)j}^i \end{pmatrix} \text{ and the duals } \begin{pmatrix} M_{(1)j}^i & M_{(2)j}^i \end{pmatrix}.$$

The invariant frames adapted to the direct decomposition

$$T_u(Osc^2M) = N_0(u) \oplus N_1(u) \oplus V_2(u) \quad \forall u \in E \quad (1)$$

will be $\mathfrak{R} = (e_{\alpha}^{(0)i}, e_{\alpha}^{(1)i}, e_{\alpha}^{(2)i})$ and the dual $\mathfrak{R}^* = (f_i^{(0)\alpha}, f_i^{(1)\alpha}, f_i^{(2)\alpha})$.

The duality conditions are:

$$\langle e_{\alpha}^{(A)i}, f_j^{(B)\alpha} \rangle = \delta_j^i \delta_B^A, \quad (A, B = 0, 1, 2).$$

In this frame the adapted basis has the representation:

$$\frac{\delta}{\delta x^i} = f^{(0)\alpha} \frac{\delta}{\delta s^{(0)\alpha}} \quad \frac{\delta}{\delta y^{(1)i}} = f^{(1)\alpha} \frac{\delta}{\delta s^{(1)\alpha}} \quad \frac{\delta}{\delta y^{(2)i}} = f^{(2)\alpha} \frac{\delta}{\delta s^{(2)\alpha}}$$

and the cobasis

$$\delta x^i = e^{(0)i} \delta s^{(0)\alpha} \quad ; \delta y^{(1)i} = e^{(1)i} \delta s^{(1)\alpha} \quad ; \delta y^{(2)i} = e^{(2)i} \delta s^{(2)\alpha} \quad (2)$$

and we have the relations:

$$\left\langle \frac{\delta}{\delta s^{(A)\alpha}}, \delta s^{(B)\beta} \right\rangle = \delta_\alpha^\beta \delta_A^B, \quad (A, B = 0, 1, 2).$$

This representation lead us to an invariant frames transformation group with the analytical expressions

$$\bar{e}^{(A)i}_\alpha = C_\alpha^A(x, y^{(1)}, y^{(2)}), e^{(A)i}_\beta \quad ; \quad f^{(B)\alpha}_j = \bar{C}_\beta^B \bar{J}^{(B)\beta}_j,$$

isomorphic with the multiplicative nonsingular matrix group

$$\begin{pmatrix} C_\beta^0 & 0 & 0 \\ 0 & C_\beta^1 & 0 \\ 0 & 0 & C_\beta^2 \end{pmatrix}.$$

A N-linear connection D has in the frame \mathfrak{R} the coefficients:

$$L_{\beta\alpha}^{0A} = f^{(A)\gamma}_m \left(\frac{\delta e^{(A)m}_\beta}{\delta s^{(0)\alpha}} + e^{(0)i}_\alpha e^{(A)j}_\beta L_{ij}^m \right), \quad (A = 0, 1, 2),$$

$$C_{\beta\alpha}^{BA} = f^{(A)\gamma}_m \left(\frac{\delta e^{(A)m}_\beta}{\delta s^{(B)\alpha}} + e^{(B)i}_\alpha e^{(A)j}_\beta C_{ij}^m \right), \quad (A = 0, 1, 2; B = 1, 2).$$

Definition 1. If the vector field $X \in \mathcal{X}(E)$ has the invariant components $X^{(A)\alpha}$ ($A = 0, 1, 2$) and we denote by $'^h$ and $'^B$ the h - and the v_B , $B = 1, 2$, covariant invariant derivative operators then

$$\begin{aligned} X^{(A)\alpha} \quad)_\beta \quad ^h &= \frac{\delta X^{(A)\alpha}}{\delta s^{(0)\beta}} + L_{\varphi\beta}^{0A} X^{(A)\varphi}, \\ X^{(A)\alpha} \quad)_\beta \quad ^B &= \frac{\delta X^{(A)\alpha}}{\delta s^{(B)\beta}} + C_{\varphi\beta}^{BA} X^{(A)\varphi}. \end{aligned} \quad (3)$$

The definition of the Lie bracket conduces us to the introduction of the non-holonomy coefficients of Vranceanu:

$$\left[\frac{\delta}{\delta s^{(A)\alpha}}, \frac{\delta}{\delta s^{(B)\beta}} \right] = \underset{(AB)}{W_{\alpha\beta}^{\gamma 0}} \frac{\delta}{\delta s^{(0)\gamma}} + \underset{(AB)}{W_{\alpha\beta}^{\gamma 1}} \frac{\delta}{\delta s^{(1)\gamma}} + \underset{(AB)}{W_{\alpha\beta}^{\gamma 2}} \frac{\delta}{\delta s^{(2)\gamma}},$$

$(A, B = 0, 1, 2, \quad A \leq B)$.

2. Einstein Equations

Let us consider a metric tensor G on Osc^2M , the invariant frames $\overline{\mathfrak{R}}$ and $\overline{\mathfrak{R}}^*$ so that the quadratic form associated to the metric has the canonical representation

$$\omega^{(A)} = (\omega^{(A)1})^2 + \dots + (\omega^{(A)i_A})^2 - (\omega^{(A)i_{A+1}})^2 + \dots - (\omega^{(A)n})^2,$$

where:

$$\omega^{(A)} = g_{ij}^{(A)} \delta y^{(A)i} \delta y^{(A)j}.$$

We introduce the Vranceanu's symbols:

$$\varepsilon_{\alpha\beta}^{(A)} = \begin{cases} \delta_{\alpha\beta} & \alpha \leq i_A, \beta \leq i_A, \\ -\delta_{\alpha\beta} & \alpha > i_A, \beta > i_A, \\ 0 & \text{in rest.} \end{cases}$$

Theorem 2.1. *The Vranceanu's symbols $\varepsilon_{\alpha\beta}^{(A)}$ represents the invariant components of the tensors $g_{ij}^{(A)}$ in the frame $\overline{\mathfrak{R}}$.*

Proposition 2.1. *The frame $\overline{\mathfrak{R}}$ defined above is pseudoorthogonal.*

Proposition 2.2. *The invariant components $\varepsilon_{\alpha\beta}^{(A)}$ of the metric tensor g of the total space E satisfy the relations:*

$$\varepsilon_{\alpha\beta\gamma}^{(A)} = 0 \quad , \quad \varepsilon_{\alpha\beta}^{(A)} \binom{(B)}{\gamma} = 0, \quad (A=0,1,2; B=1,2).$$

Proposition 2.3. *The invariant components of the canonical metrical N -linear connection $CT(N)$ satisfy the relations:*

$$\begin{aligned} L_{\alpha\gamma}^{0A} \varepsilon_{\varphi\beta}^{(A)} + L_{\beta\gamma}^{0A} \varepsilon_{\alpha\varphi}^{(A)} &= 0, \\ C_{\alpha\gamma}^{BA} \varepsilon_{\varphi\beta}^{(A)} + C_{\beta\gamma}^{BA} \varepsilon_{\alpha\varphi}^{(A)} &= 0. \end{aligned}$$

The calculus of the Ricci's tensor and the scalar curvature permit us to formulate the following result:

Theorem 2.2. *The Einstein equations have the following invariant expressions in the frame $\overline{\mathfrak{R}}$:*

$$R_{\alpha\beta} - \frac{1}{2}\varepsilon_{\alpha\beta}^{(0)}R = \kappa T_{\alpha\beta}^{(0)},$$

$$S_{\alpha\beta}^{(1)} - \frac{1}{2}\varepsilon_{\alpha\beta}^{(1)}R = \kappa T_{\alpha\beta}^{(1)}, \quad S_{\alpha\beta}^{(2)} - \frac{1}{2}\varepsilon_{\alpha\beta}^{(2)}R = \kappa T_{\alpha\beta}^{(2)},$$

$$P_{\alpha\beta}^{(B)} = (-1)^{B+1}\kappa T_{\alpha\beta}^{(0B)}, \quad P_{\alpha\beta}^{(A)} = (-1)^{B+1}\kappa T_{\alpha\beta}^{(12)}, \quad (A, B=1, 2),$$

where $T_{\alpha\beta}$ and $T_{\alpha\beta}^{(A)}$ are d -tensor fields and represent the invariant components of the energy-impulse tensor.

Theorem 2.3. *Some of the equations which give the conservation law with respect to the canonical metrical N -linear connection $CT(N)$ are given by:*

$$\left(R_{\beta}^{\alpha} - \frac{1}{2} R \delta_{\beta}^{\alpha} \right)_{\alpha} + P_{\beta}^{\alpha}{}_{\alpha}^{(1)} + P_{\beta}^{\alpha}{}_{\alpha}^{(2)} = 0,$$

$$\left(S_{\beta}^{\alpha} - \frac{1}{2} R \delta_{\beta}^{\alpha} \right)_{\alpha}^{(1)} - P_{\beta}^{\alpha}{}_{\alpha}^{(2)} + P_{\beta}^{\alpha}{}_{\alpha}^{(1)} = 0,$$

$$\left(S_{\beta}^{\alpha} - \frac{1}{2} R \delta_{\beta}^{\alpha} \right)_{\alpha}^{(2)} - P_{\beta}^{\alpha}{}_{\alpha}^{(1)} + P_{\beta}^{\alpha}{}_{\alpha}^{(2)} = 0,$$

where $R_{\beta}^{\alpha} = \varepsilon^{(0)\alpha\gamma} R_{\gamma\beta}$.

3. An example of computation for Einstein equations

In this section we compute the Einstein equations in the particular case of the generalized Lagrange space $GL^{(2)n} (M, g_{ij}(x, y^{(1)}, y^{(2)}) = e^{2\sigma(x, y^{(1)}, y^{(2)})} \gamma_{ij}(x)$

In some special cases it is not necessary to choose the frames where the quadratic forms ω have canonical shape but in this cases we must consider the restrictions $e_{\alpha}^{(0)\iota} = e_{\alpha}^{(1)\iota} = e_{\alpha}^{(2)\iota} = e_{\alpha}^{\iota}$. For the considered generalized Lagrange space let us take the canonical nonlinear connection N with the coefficients:

$$N_{(1)j}^i = \gamma_{kj}^i y^{(1)k},$$

$$N_{(2)}^i = \frac{1}{2} \left(\frac{\partial \gamma_{kj}^i}{\partial x^r} - \gamma_{rh}^i \gamma_{kj}^h \right) y^{(1)k} y^{(1)r} + \gamma_{kj}^i y^{(2)k}$$

and the Berwald connection $B\Gamma(N) = (\gamma_{jk}^i, 0, 0)$.

Theorem 3.1 *The invariant components of the Berwald connection are obtained from $B\Gamma(N)$ by a deviation induced by the invariant frames.*

Direct calculus lead us to the following expressions of the invariant components of the Berwald connection:

$$\overline{B\Gamma(N)} = \left(\gamma_{\beta\alpha}^\eta + \frac{1}{2} W_{\beta\alpha}^{\eta(00)}, \frac{1}{2} W_{\beta\alpha}^{\eta(11)}, \frac{1}{2} W_{\beta\alpha}^{\eta(22)} \right).$$

Theorem 3.2 *In invariant frames the canonical metrical N -linear connection $\overline{C\Gamma(N)}$ (with respect to the metrical tensor proposed) have the coefficients*

$$\overline{C\Gamma(N)} = \left(L_{\beta\alpha}^\gamma, \underset{(1)}{C_{\beta\alpha}^\gamma}, \underset{(2)}{C_{\beta\alpha}^\gamma} \right),$$

with

$$L_{\beta\alpha}^\gamma = \underset{0}{L_{\beta\alpha}^{\gamma*}} + \Lambda_{\beta\alpha}^\gamma,$$

$$\underset{(A)}{C_{\beta\alpha}^\gamma} + \underset{A}{\Lambda_{\beta\alpha}^\gamma} \quad (A=1,2),$$

where $\underset{(A)}{L_{\beta\alpha}^{\gamma*}}$, and $\underset{(A)}{C_{\beta\alpha}^\gamma}$ are the coefficients of Berwald connection in invariant frames

and $\underset{A}{\Lambda_{\beta\alpha}^\gamma}$ are the invariant components of the deviation tensor $\underset{A}{\Lambda_{jk}^i}$ induced by the metric, which have the expressions:

$$\underset{(B)}{\Lambda_{jk}^i} = \delta_k^i \underset{(B)}{\sigma_j} + \delta_j^i \underset{(B)}{\sigma_k} - \gamma_{jk} \underset{(B)}{\sigma^i},$$

where

$$\underset{(B)}{\sigma_j} = \frac{\delta\sigma}{\delta y^{(B)j}}, \quad \underset{(B)}{\sigma^i} = \gamma^{is} \underset{(B)}{\sigma_s}, \quad (B=1,2)$$

$$(y^{(0)i} = x^i)$$

Theorem 3.3 The deviation tensor $\Lambda_{\beta\gamma}^\alpha$ with respect to $\overline{B\Gamma(N)}$ is

$$\begin{aligned} \Lambda_{\beta\gamma}^\alpha &= \frac{1}{2} g^{\alpha\eta} (g_{\beta\eta})_{\gamma} + g_{\gamma\eta})_{\beta} - g_{\beta\gamma})_{\eta} \\ \Lambda_{\beta\gamma}^\alpha &= \frac{1}{2} g^{\alpha\eta} \left(g_{\beta\eta}^{(B)} \right)_{\gamma} + g_{\gamma\eta}^{(B)} \Big|_{\beta} - g_{\beta\gamma}^{(B)} \Big|_{\eta} \end{aligned}$$

A direct calculus of curvature tensors, Ricci's tensors and curvature scalars lead us to the following result:

Theorem 3.4 Consider the space $GL^{(2)n}$ endowed with the metric $g_{ij} = e^{2\sigma(x, y^{(1)}, y^{(2)})} \gamma_{ij}$ and the canonical metrical connection $CT(N)$. The Einsein equations in the invariant frames are:

$$\begin{aligned} r_{\alpha\beta} + R_{\alpha\beta}^* + \tilde{R}_{\alpha\beta} + \frac{1}{2} g_{\alpha\beta} (r + R^* + \tilde{R}) &= \kappa T_{\alpha\beta}^0, \\ S_{\alpha\beta}^* + \tilde{S}_{\alpha\beta} + \frac{1}{2} g_{\alpha\beta} (r + R^* + \tilde{R}) &= \kappa T_{\alpha\beta}^A, \\ P_{\alpha\beta}^* + \tilde{P}_{\alpha\beta}^* &= (-1)^B \kappa T_{\alpha\beta}^{0B}, \\ P_{\alpha\beta}^* + \tilde{P}_{\alpha\beta}^* &= (-1)^B \kappa T_{\alpha\beta}^{0B}, \quad (A, B=1, 2). \end{aligned}$$

Theorem 3.5 The conservation law with respect to the canonical metrical connection $CT(N)$ in invariant frames are:

$$\begin{aligned} \left(R_{\beta}^{\alpha} - \frac{1}{2} R \delta_{\alpha}^{\beta} \right)_{\beta} + P_{\alpha}^{\beta} \Big|_{\beta}^{(1)} + P_{\alpha}^{\beta} \Big|_{\beta}^{(2)} + P_{\alpha}^{\beta} \Big|_{\beta}^{(2)} + P_{\alpha}^{\beta} \Big|_{\beta}^{(2)} &= 0, \\ \left(S_{\alpha}^{\beta} - \frac{1}{2} R \delta_{\alpha}^{\beta} \right)_{\beta}^{(1)} - P_{\alpha}^{\beta} \Big|_{\beta}^{(1)} + P_{\alpha}^{\beta} \Big|_{\beta}^{(2)} &= 0, \\ \left(S_{\alpha}^{\beta} - \frac{1}{2} R \delta_{\alpha}^{\beta} \right)_{\beta}^{(2)} - P_{\alpha}^{\beta} \Big|_{\beta}^{(2)} - P_{\alpha}^{\beta} \Big|_{\beta}^{(2)} &= 0, \end{aligned}$$

where $' \Big|_{\beta}$ and $' \Big|_{\beta}^{(B)}$ ($B=1, 2$) are the h - and the v_B - covariant invariant derivatives with respect to $CT(N)$.

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