

# The variational problem in Lagrange spaces endowed with $(\alpha, \beta)$ -metrics

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## Abstract

In this paper we study the variational problem of Lagrange spaces with  $(\alpha, \beta)$ -metrics. The results follow the classical ones and some results of Miron R. concerning Lagrange spaces.

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**Key words:** Lagrange space,  $(\alpha, \beta)$ -metric, Euler-Lagrange equations.

## 1 Introduction

Let  $(TM, \tau, M)$  be the tangent bundle of a  $C^\infty$ -differentiable real,  $n$ -dimensional manifold  $M$ . If  $(U, \varphi)$  is a local chart on  $M$ , then the coordinates of a point  $u = (x, y) \in \tau^{-1}(U) \subset TM$  will be denoted by  $(x, y)$ . Following R.Miron [1], we have the following

**Definition 1.** a) A *differentiable Lagrangian* on  $TM$  is a mapping  $L : (x, y) \in TM \rightarrow L(x, y) \in \mathbb{R}, \forall u = (x, y) \in TM$ , which is of class  $C^\infty$  on  $\widetilde{TM} = TM \setminus \{0\}$  and is continuous on the null section of the projection  $\tau : TM \rightarrow M$ , such that

$$(1.1) \quad g_{ij}(x, y) = \frac{1}{2} \frac{\partial^2 L(x, y)}{\partial y^i \partial y^j}$$

is a  $(0, 2)$ -type symmetric  $d$ -tensor field on  $TM$ .

b) A differential Lagrangian  $L$  on  $TM$  is said to be *regular* if

$$\text{rank } ||g_{ij}(x, y)|| = n, \quad \forall (x, y) \in \widetilde{TM}.$$

We will further use its contravariant  $d$ -tensor  $g^{ij}(x, y)$  given by  $g^{ik}g_{kj} = \delta_j^i$ .

c) A *Lagrange space* is a pair  $L^n = (M, L)$  formed by a smooth real  $n$ -dimensional manifold  $M$  and a regular differentiable Lagrangian  $L$  on  $M$ , for which the  $d$ -tensor field  $g_{ij}$  from (1.1) has constant signature on  $\widetilde{TM}$ .

Let  $L : TM \rightarrow \mathbb{R}$  be a differentiable Lagrangian on the manifold  $M$ , which is not necessarily regular. A curve  $c : t \in [0, 1] \rightarrow (x^i(t)) \in U \subset M$  having the image in a domain of a chart  $U$  of  $M$ , has the extension to  $\widetilde{TM}$  given by  $c^* : t \in [0, 1] \rightarrow (x^i(t), \frac{dx^i}{dt}(t)) \in \tau^{-1}(U)$ .

The integral of action of the Lagrangian  $L$  on the curve  $c$  is given by the functional

$$(1.2) \quad I(c) = \int_0^1 L \left( x(t), \frac{dx}{dt} \right) dt.$$

Consider the curves  $c_\varepsilon : t \in [0, 1] \rightarrow (x^i(t) + \varepsilon V^i(t)) \in M$ , which have the same endpoints  $x^i(0), x^i(1)$  as the curve  $c$ ,  $V^i(0) = V^i(1) = 0$  and  $\varepsilon$  is a real number, sufficiently small in absolute value, such that  $\text{Im}c_\varepsilon \subset U$ . The extension of the curve  $c_\varepsilon$  to  $TM$  is

$$c_\varepsilon^* : t \in [0, 1] \rightarrow \left( x^i(t) + \varepsilon V^i(t), \frac{dx^i}{dt} + \varepsilon \frac{dV^i}{dt} \right) \in \tau^{-1}(U).$$

The integral of action of the Lagrangian  $L$  on the curve  $c_\varepsilon$  is

$$I(c_\varepsilon) = \int_0^1 L \left( x + \varepsilon V, \frac{dx}{dt} + \varepsilon \frac{dV}{dt} \right) dt.$$

A necessary condition for  $I(c)$  to be an extremal value of  $I(c_\varepsilon)$  is

$$\left. \frac{dI(c_\varepsilon)}{d\varepsilon} \right|_{\varepsilon=0} = 0.$$

In order that the functional  $I(c)$  be an extremal value of  $I(c_\varepsilon)$  it is necessary that  $c$  be the solution of the Euler-Lagrange equations

$$E_i(L) \equiv \frac{\partial L}{\partial x^i} - \frac{d}{dt} \left( \frac{\partial L}{\partial y^i} \right) = 0, \quad y^i = \frac{dx^i}{dt}.$$

## 2 The fundamental tensor of a Lagrange space with $(\alpha, \beta)$ -metric

We consider the functions defined on  $TM$

$$\alpha(x, y) = \sqrt{\gamma_{ij}(x)y^iy^j}, \quad \beta(x) = A_i(x)y^i,$$

where  $\gamma_{ij}(x)$  is the metric tensor of a Riemannian space  $R^n = (M, \gamma_{ij}(x))$ , and  $A_i(x)$  is a covector field.

**Definition 2.** A Lagrange space  $L^n = (M, L(x, y))$  is called *with  $(\alpha, \beta)$ -metric* if the fundamental function  $L(x, y)$  is a function  $\hat{L}$ , which depends only on  $\alpha(x, y)$  and  $\beta(x, y)$ ,

$$L = \hat{L}(\alpha(x, y), \beta(x, y)).$$

We shall use the following notations

$$\begin{aligned}\partial_i \alpha &= \frac{\partial \alpha}{\partial y^i}, & \partial_i \beta &= \frac{\partial \beta}{\partial y^i}, & \partial_i \partial_j \alpha &= \frac{\partial^2 \alpha}{\partial y^i \partial y^j}, \\ \hat{L}_\alpha &= \frac{\partial \hat{L}}{\partial \alpha}, & \hat{L}_\beta &= \frac{\partial \hat{L}}{\partial \beta}, & \hat{L}_{\alpha\alpha} &= \frac{\partial^2 \hat{L}}{\partial \alpha^2}, & \hat{L}_{\alpha\beta} &= \frac{\partial^2 \hat{L}}{\partial \alpha \partial \beta}.\end{aligned}$$

**Proposition 1.** *We have the relations*

$$\begin{aligned}\partial_i \alpha &= \alpha^{-1} y_i, & \partial_i \partial_j \alpha &= \alpha^{-1} \gamma_{ij}(x) - \alpha^{-3} y_i y_j, \\ \partial_i \beta &= A_i(x), & \partial_i \partial_j \beta &= 0,\end{aligned}$$

where  $y_i = \gamma_{ij}(x) y^j$ .

We introduce the moments of the Lagrangian  $L(x, y) = \hat{L}(\alpha(x, y), \beta(x, y))$ ,

$$p_i = \frac{1}{2} \frac{\partial L}{\partial y^i} = \frac{1}{2} \left( \hat{L}_\alpha \partial_i \alpha + \hat{L}_\beta \partial_i \beta \right)$$

and we get the following result

**Proposition 2.** *The moments of the Lagrangian  $L(x, y)$  are given by*

$$p_i = \rho y_i + \rho_1 A_i,$$

where

$$(2.3) \quad \rho = \frac{1}{2} \alpha^{-1} \hat{L}_\alpha \quad \text{and} \quad \rho_1 = \frac{1}{2} \hat{L}_\beta.$$

The two scalar functions defined in (2.3) are called the principal invariants of the Lagrange space  $L^n$ .

**Proposition 3.** *The derivatives of the principal invariants of the Lagrange space  $L^n$  are given by*

$$\partial_i \rho = \rho_{-2} y_i + \rho_{-1} A_i, \quad \partial_i \rho_1 = \rho_{-1} y_i + \rho_0 A_i,$$

where

$$\rho_{-2} = \frac{1}{2} \alpha^{-2} (\hat{L}_{\alpha\alpha} - \alpha^{-1} \hat{L}_\alpha), \quad \rho_{-1} = \frac{1}{2} \alpha^{-1} \hat{L}_{\alpha\beta}, \quad \rho_0 = \frac{1}{2} \hat{L}_{\beta\beta}.$$

**Proposition 4.** *The energy*

$$E_L = y_i \frac{\partial L}{\partial y^i} - L$$

of a Lagrangian with  $(\alpha, \beta)$ -metric is given by

$$E_L = \alpha \hat{L}_\alpha + \beta \hat{L}_\beta - \hat{L}.$$

We can determine the fundamental tensor  $g_{ij}$  of the Lagrange space with  $(\alpha, \beta)$ -metric, as follows

**Theorem 1.** *The fundamental tensor  $g_{ij}$  of the Lagrange space with  $(\alpha, \beta)$ -metric is*

$$(2.4) \quad g_{ij} = \rho \gamma_{ij} + c_i c_j,$$

where  $c_i = q_{-1} y_i + q_0 A_i$  and  $q_{-1}, q_0$  satisfy the equations

$$\rho_0 = (q_0)^2, \quad \rho_{-1} = q_0 q_{-1}, \quad \rho_{-2} = (q_{-1})^2.$$

**Theorem 2.** The reciprocal tensor  $g^{ij}$  of the fundamental tensor  $g_{ij}$  is given by

$$(2.5) \quad g^{ij} = \frac{1}{\rho} \gamma^{ij} - \frac{1}{1 + c^2} c^i c^j,$$

where  $c^i = \rho^{-1} \gamma^{ij} c_j$  and  $c^i c_i = c^2$ .

### 3 Euler-Lagrange equations in Lagrange spaces with $(\alpha, \beta)$ -metric

The Euler-Lagrange equations of the Lagrange space  $L^n$  endowed with  $(\alpha, \beta)$ -metric are

$$E_i(\hat{L}) \equiv \frac{\partial \hat{L}}{\partial x^i} - \frac{d}{dt} \left( \frac{\partial \hat{L}}{\partial y^i} \right) = 0, \quad y^i = \frac{dx^i}{dt}.$$

Considering the relations

$$\begin{aligned} \frac{\partial \hat{L}}{\partial x^i} &= \hat{L}_\alpha \frac{\partial \alpha}{\partial x^i} + \hat{L}_\beta \frac{\partial \beta}{\partial x^i}, \quad \frac{\partial \hat{L}}{\partial y^i} = \hat{L}_\alpha \frac{\partial \alpha}{\partial y^i} + \hat{L}_\beta \frac{\partial \beta}{\partial y^i} \\ \frac{d}{dt} \left( \frac{\partial \hat{L}}{\partial y^i} \right) &= \frac{d \hat{L}_\alpha}{dt} \frac{\partial \alpha}{\partial y^i} + \frac{d \hat{L}_\beta}{dt} \frac{\partial \beta}{\partial y^i} + \hat{L}_\alpha \frac{d}{dt} \left( \frac{\partial \alpha}{\partial y^i} \right) + \hat{L}_\beta \frac{d}{dt} \left( \frac{\partial \beta}{\partial y^i} \right). \end{aligned}$$

by direct calculation we infer

$$E_i(\hat{L}) = \hat{L}_\alpha E_i(\alpha) + \hat{L}_\beta E_i(\beta) - \frac{\partial \alpha}{\partial y^i} \frac{d \hat{L}_\alpha}{dt} - \frac{\partial \beta}{\partial y^i} \frac{d \hat{L}_\beta}{dt}, \quad y^i = \frac{dx^i}{dt},$$

where

$$\frac{d \hat{L}_\alpha}{dt} = \hat{L}_{\alpha\alpha} \frac{d\alpha}{dt} + \hat{L}_{\alpha\beta} \frac{d\beta}{dt}, \quad \frac{d \hat{L}_\beta}{dt} = \hat{L}_{\beta\alpha} \frac{d\alpha}{dt} + \hat{L}_{\beta\beta} \frac{d\beta}{dt}.$$

Then we get

$$\begin{aligned} E_i(\hat{L}) &= \hat{L}_\alpha E_i(\alpha) + \hat{L}_\beta E_i(\beta) - \frac{\partial \alpha}{\partial y^i} \left\{ \hat{L}_{\alpha\alpha} \frac{d\alpha}{dt} + \hat{L}_{\alpha\beta} \frac{d\beta}{dt} \right\} - \\ &\quad - \frac{\partial \beta}{\partial y^i} \left\{ \hat{L}_{\beta\alpha} \frac{d\alpha}{dt} + \hat{L}_{\beta\beta} \frac{d\beta}{dt} \right\}. \end{aligned}$$

As well, we have

$$E_i(\alpha) = \frac{1}{2\alpha} E_i(\alpha^2) + \frac{1}{2} \frac{\partial \alpha}{\partial y^i} \frac{d\alpha}{dt}, \quad E_i(\beta) = F_{ir} \frac{dx^r}{dt},$$

where

$$F_{ir} = \frac{\partial A_r}{\partial x^i} - \frac{\partial A_i}{\partial x^r}$$

is the electromagnetic tensor field. Finally we have the following relation

$$\begin{aligned} E_i(\hat{L}) = & \rho E_i(\alpha^2) + 2\rho \frac{\partial \alpha}{\partial y^i} \frac{d\alpha}{dt} + 2\rho_1 F_{ir} \frac{dx^r}{dt} - \\ & - \frac{\partial \alpha}{\partial y^i} \left\{ \hat{L}_{\alpha\alpha} \frac{d\alpha}{dt} + \hat{L}_{\alpha\beta} \frac{d\beta}{dt} \right\} - \frac{\partial \beta}{\partial y^i} \left\{ \hat{L}_{\beta\alpha} \frac{d\alpha}{dt} + \hat{L}_{\beta\beta} \frac{d\beta}{dt} \right\}. \end{aligned}$$

**Proposition 5.** *The Euler-Lagrange equation in the Lagrange space  $L^n$  endowed with  $(\alpha, \beta)$ -metric are*

$$E_i(\hat{L}) = 0, \quad y^i = \frac{dx^i}{dt}.$$

**Proposition 6.** *The following relation holds:*

$$E_i(\alpha^2) = -2\gamma_{ir} \left[ \frac{d^2 x^r}{dt^2} + \{^r_{jk}\} \frac{dx^j}{dt} \frac{dx^k}{dt} \right].$$

If we have the natural parametrization of the curve  $c : t \in [0, 1] \rightarrow (x^i(t)) \in M$  relative to the Riemannian metric  $\gamma_{ij}(x)$ , then  $\alpha \left( x, \frac{dx}{dt} \right) = 1$ . Then we get:

**Theorem 3.** *In the canonical parametrization the Euler-Lagrange equations in  $L^n$  spaces with  $(\alpha, \beta)$ -metric are*

$$\begin{aligned} (3.6) \quad E_i(\hat{L}) \equiv & -2\rho\gamma_{ir} \left[ \frac{d^2 x^r}{ds^2} + \{^r_{jk}\} \frac{dx^j}{ds} \frac{dx^k}{ds} \right] + 2\rho_1 F_{ir} \frac{dx^r}{ds} - \\ & - \frac{d\beta}{ds} \left\{ \hat{L}_{\beta\alpha} \frac{\partial \alpha}{\partial y^i} + \hat{L}_{\beta\beta} \frac{\partial \beta}{\partial y^i} \right\} = 0. \end{aligned}$$

**Proposition 7.** *If the 1-form  $\beta$  is constant on the integral curve  $c$  of the Euler-Lagrange equations, then (3.6) rewrite as the Lorentz equations of the space  $L^n$ ,*

$$\frac{d^2 x^i}{ds^2} + \{^i_{jk}\} \frac{dx^j}{ds} \frac{dx^k}{ds} = \frac{\rho_1}{\rho} F^i_r \frac{dx^r}{ds},$$

where we denoted  $F_j^i = \gamma^{ik} F_{kj}$ .

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