ORTHOGONAL ADAPTED BASIS OF $T^*(Osc^kM)$

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Dedicated to Prof. M.C. Chaki and Prof. R. Miron

Abstract

Lately a big attention has been payed on the higher order geometry. Some relevant papers are mentioned in the references. R. Miron and Gh. Atanasiu in [16], [17] studied the geometry of Osc^kM . R. Miron in [19] gave the comprehensive theory of higher order geometry and its application. Here the transformation group is slightly different from that used in [19] and it will change the geometry. The adapted basis will have different form. Such an adapted basis is constructed that $T^*_{V_0}, T^*_{V_1}, \ldots, T^*_{V_k}$ are mutually orthogonal subspaces of $T^*(Osc^kM)$ with respect to the given metric G.

AMS Subject Classification: 53B40, 53C60, 58A20.

Key words: osculator bundle, adapted base, orthogonal subspaces, d-tensor field

1. Adapted basis in $T(Osc^kM)$ and $T^*(Osc^kM)$

Here Osc^kM will be defined as a C^{∞} manifold so that the transformations of form (1.1) are allowed. It is formed as a tangent space of higher order of the base manifold M.

Let $E = Osc^k M$ be a (k+1)n dimensional C^{∞} manifold. In a local chart (U, φ) a point $u \in E$ has the coordinates:

$$(x^a, y^{1a}, y^{2a}, \dots, y^{ka}) = (y^{0a}, y^{1a}, y^{2a}, \dots, y^{ka}) = (y^{\alpha a}),$$

where $x^a = y^{0a}$ and

$$a, b, c, d, e, \ldots = 1, 2, \ldots, n, \quad \alpha, \beta, \gamma, \delta, \kappa, \ldots = 0, 1, 2, \ldots, k.$$

Editor Gr.Tsagas Proceedings of the Workshop on Global Analysis, Differential Geometry and Lie Algebras, 1998, 19-28

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The following abbreviations will be used:

$$\partial_{\alpha a} = \frac{\partial}{\partial u^{\alpha a}}, \quad \alpha = 1, 2, \dots, k, \quad \partial_a = \partial_{0a} = \frac{\partial}{\partial x^a} = \frac{\partial}{\partial u^{0a}}.$$

If in an other chart (U', φ') the point $u \in E$ has the coordinates $(x^{a'}, y^{1a'}, y^{2a'}, \dots, y^{ka'})$, then in $U \cap U'$ the allowable coordinate transformations are given by:

(1)
$$x^{a'} = x^{a'}(x^1, x^2, \dots, x^n),$$

$$y^{1a'} = (\partial_a x^{a'})y^{1a} = (\partial_{0a}y^{0a'})y^{1a},$$

$$y^{2a'} = (\partial_{0a}y^{1a'})y^{1a} + (\partial_{1a}y^{1a'})y^{2a},$$

$$y^{3a'} = (\partial_{0a}y^{2a'})y^{1a} + (\partial_{1a}y^{2a'})y^{2a} + (\partial_{2a}y^{2a'})y^{3a},$$

$$\vdots$$

$$y^{ka'} = (\partial_{0a}y^{(k-1)a})y^{1a} + (\partial_{1a}y^{(k-1)a})y^{2a} + \dots + (\partial_{(k-1)a}y^{(k-1)a})y^{ka}.$$

Theorem 1.1. The transformations of type (1.1) form a group.

A nice example of a space E can be obtained if the points $(x^a) \in M$, $\dim M = n$ are considered as the points of the curve $x^a = x^a(t)$, $t \in I$ and $y^{\alpha a}$, $\alpha = 1, 2, ..., k$ are determined by:

(2)
$$y^{\alpha a} = d_t^{\alpha} x^a, \quad d_t^{\alpha} = \frac{d^{\alpha}}{dt^{\alpha}}, \ d_t = \frac{d}{dt}.$$

If in $U \cap U'$ the equation $x^{a'} = x^{a'}(x^1(t), x^2(t), \dots, x^n(t))$ is valid, then it is easy to see that

(3)
$$y^{1a'} = d_t^1 x^{a'}, \quad y^{2a'} = d_t^2 x^{a'}, \dots, y^{ka'} = d_t^k x^{a'}$$

satisfy (1.1). In [19] $y^{\alpha a} = \frac{1}{\alpha!} d_t^{\alpha} x^a$ and it results that the structure group is different from (1.1). As from (1.2) and (1.3) it follows:

(4)
$$y^{1a'} = y^{1a'}(x, y^{1a}), \quad y^{2a'} = y^{2a'}(x, y^{1a}, y^{2a}), \dots, y^{ka'} = y^{ka'}(x, y^{1a}, \dots y^{ka})$$

and from the above equation we get (1.1).

Let us introduce the notations:

(5)
$$(0)A_a^{a'} = \partial_a x^{a'}, \quad (\alpha)A_a^{a'} = d_t^{\alpha(0)}A_a^{a'} = \frac{d^{\alpha(0)}A_a^{a'}}{dt^{\alpha}}, \quad \alpha = 1, 2, \dots, k.$$

The natural basis \bar{B}^* of $T^*(E)$ is

$$\bar{B}^* = \{dy^{0a}, dy^{1a}, \dots, dy^{ka}\}.$$

The elements of \bar{B}^* are not transformed as tensors ([19], [9]). The adapted basis B^* of $T^*(E)$ is given by

(6)
$$B^* = \{ \delta y^{0a}, \delta y^{1a}, \delta y^{2a}, \dots, \delta y^{ka} \},$$

where:

(7)
$$\delta y^{0a} = dx^{a} = dy^{0a},$$

$$\delta y^{1a} = dy^{1a} + M_{0b}^{1a} dy^{0b},$$

$$\delta y^{2a} = dy^{2a} + M_{1b}^{2a} dy^{1b} + M_{0b}^{2a} dy^{0b},$$

$$\vdots$$

$$\delta y^{ka} = dy^{ka} + M_{(k-1)b}^{ka} dy^{(k-1)b} + M_{(k-2)b}^{ka} dy^{(k-2)b} + \dots + M_{0b}^{ka} dy^{0b}.$$

Theorem 1.2. The necessary and sufficient conditions that $\delta y^{\alpha a}$ are transformed as d-tensor field, i.e.

$$\delta y^{\alpha a'} = \frac{\partial x^{a'}}{\partial x^a} \delta y^{\alpha a}, \quad \alpha = 0, 1, \dots, k$$

are given by the following equations:

$$(8) M_{\alpha b}^{(\alpha+\beta)a}(\partial_{a}x^{b'}) = M_{\alpha c'}^{(\alpha+\beta)b'}\partial_{\alpha b}y^{\alpha c'} + M_{(\alpha+1)c'}^{(\alpha+\beta)b'}\partial_{\alpha b}y^{(\alpha+1)c'} + \cdots M_{(\alpha+\beta-1)c'}^{(\alpha+\beta)b'}\partial_{\alpha b}y^{(\alpha+\beta-1)c'} + \partial_{\alpha b}y^{(\alpha+\beta)c'}, 1 < \beta, \alpha+\beta < k.$$

The proof is given in [9].

From (1.8), after some calculation, we get:

$$M_{\alpha b}^{(\alpha + \beta) a \ (0)} A_a^{b'} = \binom{\alpha}{\alpha} M_{\alpha c'}^{(\alpha + \beta) b' \ (0)} A_b^{c'} + \binom{\alpha + 1}{\alpha} M_{(\alpha + 1) c'}^{(\alpha + \beta) b' \ (1)} A_b^{c'} + \dots + \binom{\alpha + \beta - 1}{\alpha} M_{(\alpha + \beta - 1) c'}^{(\alpha + \beta - 1) b'} A_b^{c'} + \binom{\alpha + \beta}{\alpha} \binom{\beta}{\alpha} A_b^{b'}.$$

From (1.5) and (1.8) it follows:

(9)
$$M_{\alpha b}^{(\alpha+\beta)a} = M_{\alpha b}^{(\alpha+\beta)a}(x, y^1, \dots, y^\beta).$$

This equation is important when the integrability conditions are examined.

The adapted basis B^* defined by (1.6) and (1.7) is different from that introduced in [16], [17], [19]. The advantage of the present basis B^* is that the functions $M_{\alpha a}^{\beta b}$ can be determined in such a way that the elements of B^* are mutually orthogonal vectors with respect to the given nondegenerated positive definite symmetric metric tensor.

The natural basis \bar{B} of T(E) is $\bar{B} = \{\partial_{0a}, \partial_{1a}, \dots, \partial_{ka}\}$. The transformation law of its elements are given in [19].

Let us denote the adapted basis of T(E) by B, where:

(10)
$$B = \{\delta_{0a}, \delta_{1a}, \delta_{2a}, \dots, \delta_{ka}\} = \{\delta_{\alpha a}\}$$

and
$$\delta_{0a} = \partial_{0a} - N_{0a}^{1b}\partial_{1b} - N_{0a}^{2b}\partial_{2b} - \cdots - N_{0a}^{kb}\partial_{kb},$$

$$\delta_{1a} = \partial_{1a} - N_{1a}^{2b}\partial_{2b} - \cdots - N_{1a}^{kb}\partial_{kb},$$

$$\vdots$$

$$\delta_{ka} = \partial_{ka} - \partial_$$

Theorem 1.3. ([9]) The necessary and sufficient conditions that B be dual to B^* ((1.6) and (1.10)) when \bar{B} is dual to \bar{B}^* i.e.

$$<\delta_{\alpha a}\delta^{\beta b}>=\delta_{\alpha}^{\beta}\delta_{a}^{b}$$

are the following relations:

$$N_{\alpha a}^{(\alpha+\beta)b} = M_{\alpha a}^{(\alpha+\beta)b} - M_{(\alpha+1)c}^{(\alpha+\beta)b} N_{\alpha a}^{(\alpha+1)c} - M_{(\alpha+2)c}^{(\alpha+\beta)b} N_{\alpha a}^{(\alpha+\beta)b} - M_{(\alpha+\beta-1)c}^{(\alpha+\beta)b} N_{\alpha a}^{(\alpha+\beta-1)c} N_{\alpha a}^{(\alpha+\beta-1)c}$$

$$(12)$$

Theorem 1.4. ([9]) The necessary and sufficient conditions that $\delta_{\alpha a}$ with respect to (1.1) are transformed as d-tensors are the following formulae:

$$(13) N_{\alpha a'}^{(\alpha+\beta)b'}(\partial_a x^{a'}) = N_{\alpha a}^{(\alpha+\beta)c} \partial_{(\alpha+\beta)c} y^{(\alpha+\beta)b'} + N_{\alpha a}^{(\alpha+\beta-1)c} \partial_{(\alpha+\beta-1)c} y^{(\alpha+\beta)b'} + \dots + N_{\alpha a}^{(\alpha+1)c} \partial_{(\alpha+1)c} y^{(\alpha+\beta)b'} - \partial_{\alpha a} y^{(\alpha+\beta)b'}.$$

The other form of (1.13) is:

$$\begin{split} N_{\alpha a'}^{(\alpha+\beta)b'\ (0)}A_a^{a'} &= \begin{pmatrix} \alpha+\beta\\ \alpha+\beta \end{pmatrix} N_{\alpha a}^{(\alpha+\beta)c\ (0)}A_c^{b'} + \begin{pmatrix} \alpha+\beta\\ \alpha+\beta-1 \end{pmatrix} N_{\alpha a}^{(\alpha+\beta-1)c\ (1)}A_c^{b'} \\ &+ \dots + \begin{pmatrix} \alpha+\beta\\ \alpha+1 \end{pmatrix} N_{\alpha a}^{(\alpha+1)c\ (\beta-1)}A_c^{b'} - \begin{pmatrix} \alpha+\beta\\ \alpha \end{pmatrix} {}_{}^{(\beta)}A_a^{b'}. \end{split}$$

From (1.12) and (1.9) we get:

(14)
$$N_{\alpha a}^{(\alpha+\beta)b} = N_{\alpha a}^{(\alpha+\beta)b}(x, y^1, y^2, \dots, y^{\beta}).$$

Theorem 1.5. The basis vectors of \bar{B} are connected with the basis vectors of B by:

Proof. From (1.11) and (1.12) it follows (1.15).

Theorem 1.6. The basis vectors of \bar{B}^* are connected with the basis covectors of B^* by :

$$dy^{0a} = \delta y^{0a}$$

$$dy^{1a} = \delta y^{1a} - N_{0e}^{1a} \delta y^{0e},$$

$$dy^{2a} = \delta y^{2a} - N_{1e}^{2a} \delta y^{1e} - N_{0e}^{2a} \delta y^{0e},$$

$$\vdots$$

$$dy^{ka} = \delta y^{ka} - N_{(k-1)e}^{ka} \delta y^{(k-1)e} - \cdots - N_{1e}^{ka} \delta y^{1e} - N_{0e}^{ka} \delta y^{0e}.$$
Proof. From (1.7) and (1.12) it follows (1.16).

Proof. From (1.7) and (1.12) it follows (1.16)

Orthogonal adapted basis in $T^*(Osc^kM)$ 2.

Let us denote by $T_H^* = T_{V_0}^*, T_{V_1}^*, T_{V_2}^*, \dots, T_{V_k}^*$ the subspaces of $T^*(E) = T^*(Osc^kM)$ spanned by $\{\delta y^{0a}\}, \{\delta y^{1a}\}, \{\delta y^{2a}\}, \dots, \{\delta y^{ka}\}$ respectively. Then

$$T^*(E) = T^*_{V_0} \oplus T^*_{V_1} \oplus T^*_{V_2} \oplus \cdots \oplus T^*_{V_k},$$

$$dimT^*(E) = n(k+1), \quad dimT^*_{V_{\alpha}} = n, \quad \alpha = 0, 1, 2, \dots, k.$$

Let us suppose that G is a symmetric nondegenerated positive definite metric tensor and in the basis \bar{B}^* is given by:

(17)
$$G = \bar{g}_{\alpha a \beta b} dy^{\alpha a} \otimes dy^{\beta b},$$

where the summation is going over all indices. In the matrix form G can be written in the following way:

(18)
$$G = \begin{bmatrix} dy^{0a} \\ dy^{1a} \\ \vdots \\ dy^{ka} \end{bmatrix}^{T} \begin{bmatrix} \bar{g}_{0a0b} & \bar{g}_{0a1b} & \dots & \bar{g}_{0akb} \\ \bar{g}_{1a0b} & \bar{g}_{1a1b} & \dots & \bar{g}_{1akb} \\ \vdots \\ \bar{g}_{ka0b} & \bar{g}_{ka1b} & \dots & \bar{g}_{kakb} \end{bmatrix} \otimes \begin{bmatrix} dy^{0b} \\ dy^{1b} \\ \vdots \\ dy^{kb} \end{bmatrix}.$$

In the basis B^* we have:

(19)
$$G = g_{\alpha b \beta b} \delta y^{\alpha a} \otimes \delta y^{\beta b}.$$

If in (2.2) d is substituted by δ , \bar{q} by q we obtain the matrix representation of G in the basis B^* .

From (1.7) it is clear, that there are so many adapted basis B^* of $T^*(E)$ as many functions $M_{\alpha b}^{(\alpha+\beta)a}$ can be found so that the coordinate transformations of type (1.1) satisfy (1.8). In this section we shall determine such an adapted basis B^* of $T^*(E)$ in which $T_{V_0}^*, T_{V_1}^*, T_{V_2}^*, \dots, T_{V_k}^*$ are mutually orthogonal subspaces of $T^*(E)$ with respect to the given metric G ((2.1)). This condition will be satisfied if $g_{\alpha a \beta b} = 0, \forall \alpha \neq \beta$.

From (2.1), (2.3) and (1.16) we get:

(20)
$$g_{\alpha a \beta b} \delta y^{\alpha a} \otimes \delta y^{\beta b} = \bar{g}_{\gamma c \delta d} (\delta y^{\gamma c} - N_{(\gamma - 1)c}^{\gamma c} \delta_y^{(\gamma - 1)e} - \dots - N_{0e}^{\gamma c} \delta y^{0e}) \otimes (\delta y^{\delta d} - N_{(\delta - 1)e}^{\delta d} \delta_y^{(\delta - 1)e} - \dots - N_{0e}^{\delta d} \delta y^{0e}),$$
$$\alpha, \beta, \gamma, \delta = 0, 1, 2, \dots, k.$$

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For the beginning we shall take k=3. After longer calculation, using the symmetry of the metric tensor $G(g_{\alpha a \beta b} = g_{\beta b \alpha a}, \bar{g}_{\alpha a \beta b} = \bar{g}_{\beta b \alpha a})$, we obtain the coefficients of $\delta y^{\alpha a} \otimes \delta y^{\beta b}$ in the following way:

$$\begin{array}{lll} \delta y^{0a}\otimes \delta y^{0b}:g_{0a\;0b}&=&\bar{g}_{0a\;0b}-\bar{g}_{1f\;0b}N_{0a}^{3f}-\bar{g}_{2f\;0b}N_{0a}^{2f}-\bar{g}_{3f\;0b}N_{0a}^{3f}\\ &&-N_{0b}^{1d}(\bar{g}_{0a\;1d}-\bar{g}_{1f\;1d}N_{0a}^{1f}-\bar{g}_{2f\;1d}N_{0a}^{2f}-\bar{g}_{3f\;1d}N_{0a}^{3f})\\ &&-N_{0b}^{2d}(\bar{g}_{0a\;2d}-\bar{g}_{1f\;2d}N_{0a}^{1f}-\bar{g}_{2f\;2d}N_{0a}^{2f}-\bar{g}_{3f\;2d}N_{0a}^{3f})\\ &&-N_{0b}^{2d}(\bar{g}_{0a\;2d}-\bar{g}_{1f\;2d}N_{0a}^{1f}-\bar{g}_{2f\;2d}N_{0a}^{2f}-\bar{g}_{3f\;2d}N_{0a}^{3f})\\ &&-N_{0b}^{2d}(\bar{g}_{0a\;2d}-\bar{g}_{1f\;2d}N_{0a}^{1f}-\bar{g}_{2f\;3d}N_{0a}^{2f}-\bar{g}_{3f\;2d}N_{0a}^{3f}),\\ \delta y^{0a}\otimes \delta y^{1b}:g_{0a\;1b}&=&\bar{g}_{0a\;1b}-\bar{g}_{1f\;1b}N_{0a}^{1f}-\bar{g}_{2f\;1b}N_{0a}^{2f}-\bar{g}_{3f\;1b}N_{0a}^{3f}\\ &&&-N_{1b}^{2d}(\bar{g}_{0a\;2d}-\bar{g}_{1f\;2d}N_{0a}^{1f}-\bar{g}_{2f\;2d}N_{0a}^{2f}-\bar{g}_{3f\;2d}N_{0a}^{3f}),\\ \delta y^{0a}\otimes \delta y^{2b}:g_{0a\;2b}&=&\bar{g}_{0a\;2b}-\bar{g}_{1f\;2b}N_{0a}^{1f}-\bar{g}_{2f\;2b}N_{0a}^{2f}-\bar{g}_{3f\;2b}N_{0a}^{3f}\\ &&&&-N_{1b}^{3d}(\bar{g}_{0a\;3d}-\bar{g}_{1f\;3d}N_{0a}^{1f}-\bar{g}_{2f\;3d}N_{0a}^{2f}-\bar{g}_{3f\;3d}N_{0a}^{3f}),\\ \delta y^{0a}\otimes \delta y^{2b}:g_{0a\;2b}&=&\bar{g}_{0a\;2b}-\bar{g}_{1f\;2b}N_{0a}^{1f}-\bar{g}_{2f\;2b}N_{0a}^{2f}-\bar{g}_{3f\;2b}N_{0a}^{3f}\\ &&&&&&&&\\ (c)&-N_{2b}^{3d}(\bar{g}_{0a\;3d}-\bar{g}_{1f\;3d}N_{0a}^{1f}-\bar{g}_{2f\;3b}N_{0a}^{2f}-\bar{g}_{3f\;3d}N_{0a}^{3f}),\\ \delta y^{0a}\otimes \delta y^{3b}:g_{0a\;3b}&=&\bar{g}_{0a\;3b}-\bar{g}_{1f\;3b}N_{0a}^{1f}-\bar{g}_{2f\;3b}N_{0a}^{2f}-\bar{g}_{3f\;3b}N_{0a}^{3f},\\ \delta y^{1a}\otimes \delta y^{1b}:g_{1a\;1b}&=&\bar{g}_{1a\;1b}-\bar{g}_{2f\;1b}N_{1a}^{2f}-\bar{g}_{3f\;3b}N_{0a}^{3f},\\ \delta y^{1a}\otimes \delta y^{2b}:g_{1a\;2b}&=&\bar{g}_{1a\;2b}-\bar{g}_{2f\;2b}N_{1a}^{2f}-\bar{g}_{3f\;2b}N_{1a}^{3f},\\ \delta y^{1a}\otimes \delta y^{2b}:g_{1a\;2b}&=&\bar{g}_{1a\;2b}-\bar{g}_{2f\;2b}N_{1a}^{2f}-\bar{g}_{3f\;3d}N_{1a}^{3f},\\ \delta y^{2a}\otimes \delta y^{3b}:g_{1a\;3b}&=&\bar{g}_{1a\;3b}-\bar{g}_{2f\;3b}N_{1a}^{2f}-\bar{g}_{3f\;3b}N_{1a}^{3f},\\ \delta y^{2a}\otimes \delta y^{2b}:g_{2a\;2b}&=&\bar{g}_{2a\;2b}-\bar{g}_{3f\;2b}N_{1a}^{3f},\\ \delta y^{2a}\otimes \delta y^{2b}:g_{2a\;2b}&=&\bar{g}_{2a\;2b}-\bar{g}_{3f\;3b}N_{2a}^{3f},\\ (h)&-N_{2b}^{3d}(\bar{g}_{2a\;3d}-\bar{g}_{3f\;3b}N_{2a}^{3f},\\ (i)&\delta y^{2a}\otimes \delta y^{3b}:g_{2a\;3b}&=&\bar{g}_{2a\;3b}-\bar{g}_{3a\;3b}.\\ \end{array}$$

From the above equation we have:

Theorem 2.1. $T_{V_3}^*$ is orthogonal to $T_{V_2}^*(g_{2a \ 3b} = 0)$ iff

$$\bar{g}_{3f\ 3b}N_{2a}^{3f} = \bar{g}_{2a\ 3b}.$$

Proposition 2.1. $T_{V_3}^*$ is orthogonal to $T_{V_1}^*(g_{1a\ 3b}=0)$ iff

$$\bar{g}_{1a\;3b} - \bar{g}_{2f\;3b} N_{1a}^{2f} - \bar{g}_{3f\;3b} N_{1a}^{3f} = 0.$$

Proposition 2.2. If $T_{V_3}^*$ is orthogonal to $T_{V_1}^*$, then $T_{V_1}^*$ is orthogonal to $T_{V_2}^*(g_{1a\ 2b}=0)$ iff

 $\bar{g}_{1a\ 2b} - g_{2f\ 2b} N_{1a}^{2f} - \bar{g}_{3f\ 2b} N_{1a}^{3f} = 0.$

Theorem 2.2. $T_{V_3}^*$ is orthogonal to $T_{V_1}^*$ and $T_{V_1}^*$ is orthogonal to $T_{V_2}^*$ iff N_{1a}^{2f} and N_{1a}^{3f} are the solutions of the matrix equation

(23)
$$\begin{bmatrix} \bar{g}_{2b\ 2f} & \bar{g}_{2b\ 3f} \\ \bar{g}_{3b\ 2f} & \bar{g}_{3b\ 3f} \end{bmatrix} \begin{bmatrix} N_{1a}^{2f} \\ N_{1a}^{3f} \\ N_{1a}^{3f} \end{bmatrix} = \begin{bmatrix} \bar{g}_{1a\ 2b} \\ \bar{g}_{1a\ 3b} \end{bmatrix}.$$

Proposition 2.3. $T_{V_3}^*$ is orthogonal to $T_{V_0}^*$ iff

$$\bar{g}_{0a\ 3b} - \bar{g}_{1f\ 3b} N_{0a}^{1f} - g_{2f\ 3b} N_{0a}^{2f} - \bar{g}_{3f\ 3b} N_{0a}^{3f} = 0.$$

Proposition 2.4. If $T_{V_0}^*$ is orthogonal to $T_{V_3}^*$, then $T_{V_0}^*$ is orthogonal to $T_{V_2}^*$ iff

$$\bar{g}_{0a\ 2b} - g_{1f\ 2b} N_{0a}^{1f} - \bar{g}_{2f\ 2b} N_{0a}^{2f} - \bar{g}_{3f\ 2b} N_{0a}^{3f} = 0.$$

Proposition 2.5. If $T_{V_0}^*$ is orthogonal to $T_{V_3}^*$ and $T_{V_2}^*$, then $T_{V_0}^*$ is orthogonal to $T_{V_1}^*$ iff

$$\bar{g}_{0a\ 1b} - \bar{g}_{1f\ 1b} N_{0a}^{1f} - \bar{g}_{2f\ 1b} N_{0a}^{2f} - \bar{g}_{3f\ 1b} N_{0a}^{3f} = 0.$$

Theorem 2.3. $T_{V_0}^*$ is orthogonal to $T_{V_1}^*$, $T_{V_2}^*$ and $T_{V_3}^*$ iff N_{0a}^{1f} , N_{0a}^{2f} and N_{0a}^{3f} are the solutions of the following equation:

(27)
$$\begin{bmatrix} \bar{g}_{1b\ 1f} & \bar{g}_{1b\ 2f} & \bar{g}_{1b\ 3f} \\ \bar{g}_{2b\ 1f} & \bar{g}_{2b\ 2f} & \bar{g}_{2b\ 3f} \\ \bar{g}_{3b\ 1f} & \bar{g}_{3b\ 2f} & \bar{g}_{3b\ 3f} \end{bmatrix} \begin{bmatrix} N_{0a}^{1f} \\ N_{0a}^{2f} \\ N_{0a}^{3f} \\ N_{0a}^{3f} \end{bmatrix} = \begin{bmatrix} \bar{g}_{0a\ 1b} \\ \bar{g}_{0a\ 2b} \\ \bar{g}_{0a\ 3b} \end{bmatrix}.$$

Theorem 2.4. The necessary and sufficient conditions that the subspaces $T_{V_0}^*$, $T_{V_1}^*$, $T_{V_2}^*$ and $T_{V_3}^*$ of $T^*(Osc^3M)$ formed by $\{\delta y^{0a}\}$, $\{\delta y^{1a}\}$, $\{\delta y^{2a}\}$ and $\{\delta y^{3a}\}$ respectively are mutually orthogonal with respect to the given metric G (given by (2.2)) are the equations (2.5), (2.7) and (2.11).

Theorem 2.5. When $T_{V_0}^*$, $T_{V_1}^*$, $T_{V_2}^*$ and $T_{V_3}^*$ are mutually orthogonal subspaces of $T^*(Osc^3M)$, with respect to the metric G, then:

$$g_{0a\ 0b} = \bar{g}_{0a\ 0b} - \bar{g}_{1f\ 0b} N_{0a}^{1f} \bar{g}_{2f\ 0b} N_{0a}^{2f} - \bar{g}_{3f\ 0b} N_{0a}^{3f},$$

$$g_{1a\ 1b} = \bar{g}_{1a\ 1b} - \bar{g}_{2f\ 1b} N_{1a}^{2f} - \bar{g}_{3f\ 1b} N_{1a}^{3f},$$

$$g_{2a\ 2b} = \bar{g}_{2a\ 2b} - \bar{g}_{3f\ 2b} N_{2a}^{3f},$$

$$g_{3a\ 3b} = \bar{g}_{3a\ 3b}.$$

Let us introduce the notations:

(29)
$$\bar{G}_{\beta k} = \begin{bmatrix} \bar{g}_{\beta b \beta f} & \cdots & \bar{g}_{\beta b k f} \\ \vdots & & & \\ \bar{g}_{k b \beta f} & \cdots & \bar{g}_{k b k f} \end{bmatrix}, \\
\bar{G}^{\beta k} = (G_{\beta k})^{-1} = \begin{bmatrix} \bar{g}^{\beta b \beta f} & \cdots & \bar{g}^{\beta b k f} \\ \bar{g}^{k b \beta f} & \cdots & \bar{g}^{k b k f} \end{bmatrix}, \\
\bar{G}_{\alpha, \beta k} = \begin{bmatrix} \bar{g}_{\alpha a \beta b} \\ \vdots \\ \bar{g}_{\alpha a k b} \end{bmatrix}.$$

From (2.13) it is clear that $\bar{G}_{\beta k}$ and $\bar{G}^{\beta k}$ are matrices of type $(k-\beta+1)\times(k-\beta+1)$ and $\bar{G}_{\alpha,\beta k}$ is a matrix of type $(k-\beta+1)\times 1$. The elements of all three matrices are submatrices of type $n\times n$.

As in all propositions and theorems of this section it was supposed that k = 3, so equations (2.5), (2.7) and (2.11) can be written in the form:

$$N_{2a}^{3b} = \bar{g}^{3b} \,^{3e} g_{2a} \,_{3b} = \bar{G}^{33} \bar{G}_{2,33},$$

$$\begin{bmatrix} N_{1a}^{2e} \\ N_{1a}^{3e} \end{bmatrix} = \begin{bmatrix} \bar{g}^{2e} \, ^{2b} & \bar{g}^{2e} \, ^{3b} \\ \bar{g}^{3e} \, ^{2b} & \bar{g}^{3e} \, ^{3b} \end{bmatrix} \begin{bmatrix} \bar{g}_{1a} \, ^{2b} \\ \bar{g}_{1a} \, ^{3b} \end{bmatrix} = \bar{G}^{23} \bar{G}_{1,23},$$

$$\begin{bmatrix} N_{0a}^{1e} \\ N_{0a}^{2e} \\ N_{0a}^{3e} \\ N_{0a}^{3e} \end{bmatrix} = \begin{bmatrix} \bar{g}^{1e \ 1b} & \bar{g}^{1e \ 2b} & \bar{g}^{1e \ 3b} \\ \bar{g}^{2e \ 1b} & \bar{g}^{2e \ 2b} & \bar{g}^{2e \ 3b} \\ \bar{g}^{3e \ 1b} & \bar{g}^{3e \ 2b} & \bar{g}^{3e \ 3b} \end{bmatrix} \begin{bmatrix} \bar{g}_{0a \ 1b} \\ \bar{g}_{0a \ 2b} \\ \bar{g}_{0a \ 3b} \end{bmatrix} = \bar{G}^{13} \bar{G}_{0,13}.$$

The matrices on the right hand side in (2.14), (2.15) and (2.16) are the correspondent inverse matrices which appears in (2.5), (2.7) and (2.11).

Now we have:

Theorem 2.4'. The necessary and sufficient conditions that the subspaces $T_{V_0}^*$, $T_{V_1}^*$, $T_{V_2}^*$ and $T_{V_3}^*$ of $T^*(Osc^3M)$ formed by $\{\delta y^{0a}\}$, $\{\delta y^{1a}\}$, $\{\delta y^{2a}\}$ and $\{\delta y^{3a}\}$ respectively be mutually orthogonal with respect to the given metric G are the equations (2.14), (2.15) and (2.16).

The main result is the following theorem:

Theorem 2.6. If in $T^*(Osc^kM)$ the metric tensor G is given by (2.2), then there exists one and only one adapted basis $\{\delta y^{0a}, \delta y^{1a}, \ldots, \delta y^{ka}\}$ such that the subspaces $T^*_{V_0}, T^*_{V_1}, \ldots, T^*_{V_k}$ of $T^*(Osc^kM)$ are mutually orthogonal. The vectors of such unique base are determined by (1.16) and the coefficients N are given by:

$$\begin{bmatrix} N_{\alpha-1}^{\alpha e} \\ N_{\alpha-1}^{(\alpha+1)e} \\ N_{\alpha-1}^{ke} \end{bmatrix} = \bar{G}^{\alpha k} \bar{G}_{\alpha-1,\alpha k} \quad \alpha = 1, 2, \dots, k.$$

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