

NOTES ON SUBGEODESIC AND GEODESIC MAPPINGS OF RIEMANNIAN MANIFOLDS

Iulia - Elena Hirićă

Abstract

The present paper deals with subgeodesically and geodesically related Riemannian spaces. In the third section we study a subprojective transformation. In the last section, using properties of conharmonic, concircular transformations and tensors satisfying some pseudo-symmetry conditions we study subgeodesically related manifolds.

Key words: subgeodesic mappings, pseudo-symmetric spaces, conharmonic transformations, concircular transformations.

Mathematics Subject Classification: 53B20, 53C25.

1 Introduction

Let $V_n = (M, g)$ be a Riemannian manifold. It is said to be pseudo-symmetric ([6]) if

(1.1) at every point of M the tensors $R \cdot R$ and $Q(g, R)$ are linearly dependent.

It is clear that any semi-symmetric manifold $R \cdot R = 0$ ([18]) is pseudo-symmetric. These notions arose during the study of totally umbilical submanifolds of semi-symmetric manifolds ([1]) as well as during the consideration of geodesic mappings ([3], [7], [8]).

In this paper we shall continue the study in this direction, considering subgeodesic mappings ([13], [16]) which are a generalization of the geodesic mappings. We determine some properties of ξ^i -subgeodesically related spaces, the tensor of correspondence being $-g$, with conharmonically semi-symmetric spaces or pseudo-symmetric spaces. We study also the Weyl conformal curvature tensor produced by some subprojective transformations.

Editor Gr.Tsagas *Proceedings of the Workshop on Global Analysis, Differential Geometry and Lie Algebras, 1996, 205-212*

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2 Preliminaries

Let (M, g) be a Riemannian manifold, covered by a system of coordinates neighbourhoods. We denote by g_{ij} , R_{ijkl} , R_{ij} and K the components of the Riemannian metric, the Riemannian curvature tensor, the Ricci tensor and the scalar curvature, respectively. We define the tensors $R(1)$, $R \cdot R$ and $Q(g, R)$ by the formulae

$$(2.1) \quad R(1)_{ijkl} = \frac{K}{n(n-1)}(g_{ik}g_{jh} - g_{ij}g_{kh}),$$

$$(2.2) \quad (R \cdot R)_{hijklm} = -R_{hlm}^s R_{sijk} - R_{ilm}^s R_{hsjk} - R_{jlm}^s R_{hisl} - R_{klm}^s R_{hij s},$$

$$(2.3) \quad \begin{aligned} Q(g, R)_{hijkl} = & g_{mh} R_{lij k} - g_{hl} R_{mijk} + g_{mi} R_{hljk} - \\ & - g_{il} R_{hmjk} + g_{jm} R_{lik} - g_{jl} R_{himk} + g_{km} R_{hijl} - g_{kl} R_{hijm}, \end{aligned}$$

where $R_{hlm}^s = g^{sr} R_{rhlm}$.

The condition (1.1) holds on $U = \{x \in M \mid R \neq R(1) \text{ at } x\}$ if and only if

$$(*) \quad R \cdot R = LQ(g, R)$$

is satisfied on U , where L is a function.

We can also define the tensors $R \cdot g$ and $Q(g, B)$, B being a symmetric tensor field of type $(0, 2)$.

$$(2.4) \quad (R \cdot g)_{jkri} = g_{jk,ir} - g_{jk,ri},$$

$$(2.5) \quad Q(g, B)_{jkri} = g_{ij} B_{kr} - g_{rj} B_{ki} - g_{ki} B_{jr} - g_{rk} B_{ji},$$

where the comma denotes covariant differentiation with respect to the metric g .

A $(0, 4)$ -tensor field T on M is said to be a generalized curvature tensor if

$$(2.6) \quad T_{ijkl} + T_{iklj} + T_{iljk} = 0,$$

$$(2.7) \quad T_{ijkl} = -T_{jikl},$$

$$(2.8) \quad T_{ijkl} = T_{klij}.$$

Further, we define the tensor $T \cdot T$ analogous with $R \cdot R$ and the Weyl conformal curvature tensor $W(T)$ associated to T by the formula ([5])

$$(2.9) \quad \begin{aligned} (W(T))_{ijkl} = & T_{ijlk} - \frac{1}{n-2}(g_{ik}T_{jl} - g_{il}T_{jk} - g_{jk}T_{il} - \\ & - g_{jl}T_{ik}) + \frac{K(T)}{(n-1)(n-2)}(g_{ik}g_{jl} - g_{il}g_{kj}), \end{aligned}$$

where $T_{hlm}^s = g^{rs} T_{rhlm}$, $T_{il} = T_{ir}^r$ and $K(T)$ is the scalar curvature associated to T .

The conharmonic curvature tensor on $V_n = (M, g)$ is defined by ([11])

$$(2.10) \quad C_{ijkl} = R_{ijkl} - \frac{1}{n-2}(g_{ik}R_{jl} - g_{il}R_{jk} - g_{jk}R_{il} + g_{jl}R_{ik}).$$

A manifold is said to be conharmonically semi-symmetric if $R \cdot C = 0$, the tensor $R \cdot C$ being defined by analogy with (2.2), replacing the Riemann tensor with the conharmonic tensor C .

Let A_n be a space with an affine, symmetric connection Γ_{jk}^i and ξ^i be the components of a vector field on A_n . We can associate the differential system of equations

$$\frac{d^2 x^i}{dt^2} + \Gamma_{jk}^i \frac{dx^j}{dt} \frac{dx^k}{dt} = a \frac{dx^i}{dt} + b \xi^i,$$

a and b being functions of t , which defines the ξ^i -subgeodesics.

Let $A_n(\Gamma_{jk}^i)$ and $\bar{A}_n(\bar{\Gamma}_{jk}^i)$ be two spaces. K. Yano introduced ([23]) the subprojective transformation of connections, which preserves the ξ^i -subgeodesics

$$(2.11) \quad \bar{\Gamma}_{jk}^i = \Gamma_{jk}^i + \delta_j^i \psi_k + \delta_k^i \psi_j + \phi_{jk} \xi^i,$$

where ψ_i and ϕ_{jk} are the components of a 1-form and of a symmetric tensor field of type $(0, 2)$, respectively.

We say that two Riemannian spaces $V_n = (M, g)$ and $\bar{V}_n = (M, \bar{g})$ are ξ^i -subgeodesically related ([13], [16]), the tensor of correspondence being $-g_{ij}$, if we have Yano formulae

$$(2.12) \quad \overline{\left| \begin{matrix} i \\ jk \end{matrix} \right|} = \left| \begin{matrix} i \\ jk \end{matrix} \right| + \delta_j^i \psi_k + \delta_k^i \psi_j - g_{kj} \xi^i,$$

where $\overline{\left| \begin{matrix} i \\ jk \end{matrix} \right|}$, $\left| \begin{matrix} i \\ jk \end{matrix} \right|$ are the Christoffel symbols for \bar{V}_n and V_n , respectively. ψ_i and ξ^i are the components of a 1-form and of a vector field, respectively.

This is equivalent with the existence of a diffeomorphism f between these two spaces which maps ξ^i -subgeodesics onto ξ^i -subgeodesics. f is called the subgeodesic mapping.

From (2.12) we have

$$(n+1)\psi_k - \xi_k = \frac{\partial}{\partial x^k} \ln \sqrt{\left| \frac{\det(\bar{g}_{ij})}{\det(g_{ij})} \right|},$$

where $\xi_k = g_{ks} \xi^s$. There exist the functions $\psi(x^1, \dots, x^n), \xi(x^1, \dots, x^n)$ ([13]) such that $\psi_k = \frac{\partial \psi}{\partial x^k}$ and $\xi_k = \frac{\partial \xi}{\partial x^k}$.

Spaces V_n and \bar{V}_n are non-trivial ξ^i -subgeodesically related, the tensor of correspondence being $-g_{ij}$, if the components $\psi_i - \xi_i$ are non-zero.

If $\xi^i = 0$ then the Yano formulae become the Weyl formulae and V_n and \bar{V}_n are geodesically related spaces ([2], [15], [21]).

3 Subprojective transformations

On a Riemannian manifold $V_n = (M, g)$ we consider the linear connection

$$(3.1) \quad \overline{\Gamma}_{jk}^i = \left| \begin{array}{c} i \\ j \quad k \end{array} \right| + \delta_j^i \Psi_k + \delta_k^i \Psi_j - g_{kj} \xi^i.$$

Using (3.1) we have

$$(3.2) \quad T_{ijkl} = R_{ijkl} - g_{ij}(A_{kl} - A_{lk}) + g_{il}A_{jk} - g_{ik}A_{jl} + \\ + g_{jk}B_{il} - g_{jl}B_{ik} - (g_{ik}g_{jl} - g_{il}g_{jk})\Psi_s \xi^s,$$

where $T_{ijkl} = g_{is} \overline{R}_{jkl}^s$, $A_{ij} = \Psi_{i,j} - \Psi_i \Psi_j$, $B_{ij} = \xi_{i,j} - \xi_i \xi_j$ and the comma denotes covariant differentiation with respect to g .

Theorem 3.1. *T , defined by (3.2), is a generalized curvature tensor if and only if the relations*

$$(3.3) \quad B_{ij} = A_{ij} - \frac{1}{n} g_{ij} [T_r(A) - T_r(B)], \quad A_{ij} = A_{ji}, \quad n \geq 3,$$

hold.

Proof. The condition (2.6) is identically satisfied. The relation (2.7) is equivalent to

$$(3.4) \quad g_{il}(A_{jk} - B_{jk}) + g_{jl}(A_{ik} - B_{ik}) - g_{jk}(A_{il} - B_{il}) - \\ - g_{ik}(A_{jl} - B_{jl}) = 2g_{ij}(A_{kl} - A_{lk}).$$

Transvecting with g^{ij} and summing, (3.4) leads to

$$(3.5) \quad B_{kl} - B_{lk} = (n+1)(A_{kl} - A_{lk}).$$

The relation (2.8) is equivalent with

$$(3.6) \quad g_{ij}(A_{kl} - A_{lk}) + g_{ik}(A_{jl} - A_{lj}) + g_{jl}(B_{ik} - B_{ki}) - \\ - g_{il}(A_{jk} - B_{kj}) + g_{jk}(A_{li} - B_{il}) = g_{kl}(A_{ij} - A_{ji}).$$

Transvecting with g^{ik} in (3.6), we have

$$(3.7) \quad B_{jl} = \frac{n+2}{n} A_{jl} - \frac{2}{n} A_{lj} - \frac{1}{n} g_{jl} [T_r(A) - T_r(B)].$$

Using (3.6) and (3.7), we obtain (3.2), for $n \geq 3$.

If the conditions (3.2) are satisfied, then T verifies the relations (2.6), (2.7) and (2.8).

Theorem 3.2. *If T , defined by (3.2), is a generalized curvature tensor on $V_n = (M, g)$, $n \geq 3$, then the Weyl curvature tensor of V_n coincides with the Weyl curvature tensor associated with T .*

Proof. Substituting the relation (3.3) in (3.2) we get

$$(3.8) \quad T_{ijkl} = R_{ijkl} + g_{il}A_{jk} - g_{ik}A_{jl} + g_{jk}A_{il} - g_{jl}A_{ik} + (g_{ik}g_{jl} - g_{il}g_{jk}) \left[\frac{1}{n}(T_r(A) - T_r(B)) - \Psi_s \xi^s \right].$$

In view of (3.8), we obtain

$$(3.9) \quad T_{jl} = R_{jl} - (n-2)A_{jl} - g_{jl}T_r(A) + (n-1)g_{jl} \left[\frac{1}{n}(T_r(A) - T_r(B)) - \Psi_s \xi^s \right].$$

Using (3.9) we have

$$(3.10) \quad K(T) = K - (n-2)T_r(A) - nT_r(A) + n(n-1) \left[\frac{1}{n}(T_r(A) - T_r(B)) - \Psi_s \xi^s \right].$$

The relations (3.9) and (3.10) lead to

$$(3.11) \quad A_{jl} = \frac{R_{jl} - T_{jl}}{n-2} + \frac{K(T) - K}{n(n-2)}g_{jl} + \frac{1}{n}T_r(A)g_{jl}.$$

Substituting (3.10) and (3.11) in (3.8), we obtain $W = W(T)$.

Proposition 3.1. Let T , defined by (3.2), be a generalized curvature tensor on the Riemannian manifold $V_n = (M, g)$, $n \geq 3$. If the Weyl conformal curvature tensor W verifies the relation

$$(3.12) \quad \tau(X)W(Y, Z) + \tau(Y)W(Z, X) + \tau(Z)W(X, Y) = 0,$$

where τ is a 1-form, then the condition $W(T) \cdot W(T) = 0$ holds at x , where $\tau \neq 0$ at x , $W(T)$ being the Weyl conformal curvature tensor associated with T .

Proof. If the relation (3.12) is satisfied, then the condition $W \cdot W = 0$ holds at x , where $\tau \neq 0$ at x ([5]).

Using the previous theorem we obtain our assertion.

4 Subgeodesically and geodesically related Riemannian spaces

Theorem 4.1. Let $V_n = (M, g)$ and $\bar{V}_n = (M, \bar{g})$, $n \geq 3$, be two ξ^i -subgeodesically related Riemannian spaces, the tensor of correspondence being $-g$, such that $\xi_{hk} = 0$, where $\xi_{hk} = \xi_{h,k} - \xi_h \xi_k + \frac{1}{2} \xi_i \xi^i g_{hk}$.

If \bar{V}_n is a conharmonically semi-symmetric space and with irreducible curvature tensor, then V_n is an Einstein space.

Proof. Let $\tilde{V}_n = (M, \tilde{g} = e^{2\xi}g)$. V_n and \tilde{V}_n being conformally related, we have

$$(4.1) \quad \left| \widetilde{\begin{matrix} i \\ j \ k \end{matrix}} \right| = \left| \begin{matrix} i \\ j \ k \end{matrix} \right| + \delta_j^i \xi_k + \delta_k^i \xi_j - g_{jk} \xi^i.$$

Using (2.12) and (4.1) we get $\left| \overline{\begin{matrix} i \\ j \ k \end{matrix}} \right| = \left| \widetilde{\begin{matrix} i \\ j \ k \end{matrix}} \right| + \delta_j^i \omega_k + \delta_k^i \omega_j$, where $\omega_i = \Psi_i - \xi_i$.

Hence \overline{V}_n and \tilde{V}_n are geodesically related spaces. In the same way as in the proofs of the theorems 1-3 of [10], \tilde{V}_n is an Einstein space.

A necessary and sufficient condition that an Einstein space be transformed into an Einstein space by a conharmonic transformation is that $\xi_{ij} = 0$ ([11]).

Because $\xi_{ij} = 0$, $g \rightarrow \tilde{g} = e^{2\xi}g$ is a conharmonic transformation i.e. $\xi_p^p = g^{ij} \xi_{ij} = 0$. Applying the previous property, V_n will be an Einstein space.

Theorem 4.2. Let $V_n = (M, g)$ and $\overline{V}_n = (M, \overline{g})$, $n \geq 3$, be two ξ^i -subgeodesically related Riemannian spaces, the tensor of correspondence being $-g$, such that $B = \frac{1}{n} T_r(B)g$, where B is a tensor field of type $(0, 2)$, having the components $B_{rs} = \xi_{r,s} - \xi_r \xi_s$.

If V_n is a pseudo-symmetric space, then \overline{V}_n is a pseudo-symmetric space.

Proof. If $B = \frac{1}{n} T_r(B)g$, then $g \rightarrow \tilde{g} = e^{2\xi}g$ is a concircular transformation. If V_n is a pseudo-symmetric space, then \tilde{V}_n is a pseudo-symmetric space ([4]).

\overline{V}_n and \tilde{V}_n being geodesically related manifolds, the space \overline{V}_n will be pseudo-symmetric ([3], [12]).

Theorem 4.3. Let $V_n = (M, g)$ and $\overline{V}_n = (M, \overline{g})$, $n \geq 3$, be two non-trivial ξ^i -subgeodesically related Riemannian spaces, the tensor of correspondence being $-g$.

Then $\overline{R} \cdot g = Q(g, F)$, where F is a symmetric tensor field of type $(0, 2)$, having the components

$$(4.2) \quad \overline{F}_{ij} = \xi_{i;j} - \Psi_{i;j} - (\xi_i - \Psi_i)(\xi_j - \Psi_j),$$

where ";" denotes covariant differentiation with respect to the metric \overline{g} .

Proof. Using (2.12) we get

$$\begin{aligned} g_{jk;ir} &= -2\Psi_{i;r}g_{jk} - (\Psi_{j;r} - \xi_{j;r})g_{ik} - (\Psi_{k;r} - \xi_{k;r})g_{ij} - \\ &\quad - 2\Psi_i[-2\Psi_r g_{jk} - (\Psi_j - \xi_j)g_{rk} - (\Psi_k - \xi_k)g_{rj}] - \\ &\quad - (\Psi_j - \xi_j)[-2\Psi_r g_{ik} - (\Psi_i - \xi_i)g_{rk} - (\Psi_k - \xi_k)g_{ir}] - \\ &\quad - (\Psi_k - \xi_k)[-2\Psi_r g_{ij} - (\Psi_j - \xi_j)g_{ri} - (\Psi_i - \xi_i)g_{rj}]. \end{aligned}$$

Hence

$$(\overline{R} \cdot g)_{jkri} = g_{jk;ir} - g_{jk;ri} = Q(g, F)_{jkri},$$

where $\overline{F}_{ij} = \xi_{i;j} - \Psi_{i;j} - (\xi_i - \Psi_i)(\xi_j - \Psi_j)$.

Proposition 4.1. Let $V_n = (M, g)$ and $\bar{V}_n = (M, \bar{g})$, $n \geq 3$, be two non-trivial ξ^i -subgeodesically related Riemannian spaces, the tensor of correspondence being $-g$.

We suppose that \bar{V}_n is a pseudo-symmetric space and the function \bar{L} , which satisfies the condition (*) on \bar{U} , is constant.

If $F_{ij} = fg_{ij} + h\bar{g}_{ij}$, where $f, h \in \mathcal{F}(M)$, h being non-nulle and F_{ij} defined by (4.2), then the relation :

$$(4.3) \quad (\bar{L} + h) \left[g - \frac{1}{n} (\bar{g}^{ij} g_{ij}) \bar{g} \right] = 0$$

holds on \bar{U} .

Proof. We will apply the property ([7]) :

Let A and D be symmetric tensors of type $(0,2)$ on a pseudo-symmetric manifold $V_n = (M, g)$, having a constant function L , which satisfies the condition (*) on U .

If the relation $R \cdot A = Q(g, D)$ holds on U , then the relation $E - \frac{1}{n} T_r(E)g = 0$ is satisfied on U , where $E = D - LA$.

Because $F = fg + h\bar{g}$, using the theorem 4.3, we have $\bar{R} \cdot g = Q(\bar{g}, -hg)$.

In this case, the tensor $E = -hg - \bar{L}g$ satisfies on \bar{U} the relation

$$E - \frac{1}{n} (\bar{g}^{ij} E_{ij}) \bar{g} = 0.$$

This condition is equivalent with (4.3).

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Author's address:

Iulia - Elena Hirićă
 University of Bucharest
 Faculty of Mathematics
 Department of Geometry
 14 Academiei St., code 70109
 e-mail: ihirica@geo.math.unibuc.ro