

# ON THE LAGRANGE SPACES WITH $(\alpha, \beta)$ -METRICS

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## Abstract

One considers a Lagrange space with the fundamental function (1.1) and one determines the fundamental tensor field, canonical nonlinear connection and canonical metrical linear N-connection.

## 1. Introduction

Let  $M$  be a real,  $n$ -dimensional  $C^\infty$  manifold and  $\tau : TM \rightarrow M$  its tangent bundle. We shall consider, together with a Finsler space  $F^n = (M, F(x, y))$ , a covector field  $A_i(x)$  defined on  $M$ , or on open set of  $M$ , and "the electromagnetic" 1-form  $\beta(x, y) = A_i(x)y^i$ . So, we can define the real function on  $TM$ :

$$(1.1) \quad L(x, y) = \alpha^2(x, y) + a\beta(x, y) + b\beta^2(x, y), \quad \forall (x, y) \in TM,$$

where  $\alpha(x, y) = F(x, y)$ .

This function is  $C^\infty$ -differentiable on  $\widetilde{TM} = TM - \{O\}$  and continuous on the null section  $O : M \rightarrow TM$ . Obviously  $L(x, y)$  is not homogeneous with respect to  $(y^i)$ . We prove that  $L(x, y)$  is a regular Lagrangian and that  $L^n = (M, L(x, y))$  is a Lagrange space [2]. We study the Lagrange space  $L^n$  with the fundamental function (1.1) and determine the fundamental tensor field, canonical nonlinear connection and canonical metrical linear  $N$ -connection.

## 2. The Lagrange space $L^n = (M, L)$

According to (1.1) the fundamental tensor field of the Lagrange space  $L^n$ ,  $g_{ij} = \frac{1}{2} \frac{\partial^2 L}{\partial y^i \partial y^j}$ , is given by:

**Proposition 2.1** *The fundamental tensor field of the Lagrange space  $L^n = (M, L)$  is:*

$$(2.1) \quad g_{ij}(x, y) = \gamma_{ij}(x, y) + bA_i(x)A_j(x),$$

where  $\gamma_{ij} = \frac{1}{2} \frac{\partial^2 F^2}{\partial y^i \partial y^j}$  is the fundamental tensor field of the Finsler space  $F^n$ .

In order to deduce that  $g_{ij}$ , from (2.1), is nonsingular we distinguish two cases:  $b > 0$  and  $b < 0$ . In this respect we reproduce the following Lemmas from the paper [2].

**Lemma 2.2** *Let  $\|A_{ij}\|$ ,  $(i, j = 1, \dots, n)$  a real nonsingular matrix, having  $\|A_{ij}\|^{-1} = \|A^{ij}\|$ . Then the matrix  $\|B_{ij}\|$  with the elements  $B_{ij} = A_{ij} + c_i c_j$ , such that  $1 + c^2 \neq 0$ ,  $c^2 = A^{ij} c_i c_j$ , is nonsingular. Its determinant is  $\det\|B_{ij}\| = (1 + c^2)\det\|A_{ij}\|$  and  $\|B_{ij}\|^{-1}$  has the elements  $B^{ij} = A^{ij} - \frac{1}{1 + c^2} c^i c^j$ , ( $c^i = A^{ij} c_j$ ).*

**Lemma 2.3** *If  $\|A_{ij}\|$ ,  $(i, j = 1, \dots, n)$  is a real nonsingular matrix, having  $A^{ij}$  as elements of its inverse and  $d_i$ ,  $(i = 1, \dots, n)$  are real numbers for which  $1 - d^2 \neq 0$ ,  $d^2 = A^{ij} d_i d_j$ , then the matrix with the elements  $B_{ij} = A_{ij} - d_i d_j$  is nonsingular. It has the determinant  $\det\|B_{ij}\| = (1 - d^2)\det\|A_{ij}\|$  and its inverse has the elements  $B^{ij} = A^{ij} + \frac{1}{1 - d^2} d^i d^j$ ,  $d^i = A^{ij} d_j$ .*

Applying these Lemmas we obtain:

**Theorem 2.1** <sup>10</sup> *The d-tensor field  $g_{ij}$  has the following properties:*

$$(2.2) \quad \begin{aligned} &\text{If } b > 0, \text{ then } \det\|g_{ij}\| = (1 + c^2)\det\|\gamma_{ij}\|, \\ &\text{where } c^2 = b\gamma^{ij}(x, y)A_i(x)A_j(x). \\ &\text{If } b < 0, \text{ then } \det\|g_{ij}\| = (1 - d^2)\det\|\gamma_{ij}\|, \\ &\text{where } d^2 = -b\gamma^{ij}(x, y)A_i(x)A_j(x). \end{aligned}$$

<sup>20</sup> *The contravariant tensor of  $(g_{ij})$  is as follows:*

$$(2.3) \quad \begin{aligned} &\text{If } b > 0, \text{ then } g^{ij}(x, y) = \gamma^{ij}(x, y) - \frac{1}{1 + c^2} A^i(x, y)A^j(x, y), \\ &\text{where } A^i(x, y) = \sqrt{b}\gamma^{ij}(x, y)A_j(x). \\ &\text{If } b < 0, \text{ then } g^{ij}(x, y) = \gamma^{ij}(x, y) + \frac{1}{1 - d^2} A^i(x, y)A^j(x, y), \\ &\text{where } A^i(x, y) = \sqrt{-b}\gamma^{ij}(x, y)A_j(x). \end{aligned}$$

Now, we can state:

**Theorem 2.2** *The pair  $L^n = (M, L(x, y))$ , where  $L(x, y)$  is given by (1.1), is a Lagrange space.*

**Remark.** The classical case is obtained when  $\gamma_{ij}(x, y) = \gamma_{ij}(x)$ , (Lorentz metric).

The space  $L^n$  is called the Lagrange space of generalized electrodynamics and  $F^n$  the associated Finsler space to  $L^n$ .

### 3. Variational problem

Let  $c : [0, 1] \rightarrow M$  be a smooth curve in  $M$  expressed in a local chart  $(U, \varphi)$  on the base manifold  $M$  by  $x^i = x^i(t), t \in [0, 1], \text{Im } c \subset U$ . The length of  $c$  in the Lagrange

space  $L^n$  is

$$(3.1) \quad I(c) = \int_0^1 L(x, \dot{x}) dt.$$

The variational problem concerning  $I(c)$ , leads to the Euler-Lagrange equation:

$$(3.2) \quad \begin{cases} \frac{d}{dt} \left\{ \frac{\partial L}{\partial y^i} \right\} - \frac{\partial L}{\partial x^i} = 0 \\ y^i = \frac{dx^i}{dt}. \end{cases}$$

We denote the electromagnetic tensor field, determined by the covector field  $A_i(x)$ , by

$$(3.3) \quad F_{ij}(x) = \frac{\partial A_j}{\partial x^i} - \frac{\partial A_i}{\partial x^j}$$

and consider its mixed form

$$(3.4) \quad F_j^i(x, y) = \gamma^{ik}(x, y) F_{kj}(x).$$

After usual calculation we get:

**Theorem 3.1** *The Euler-Lagrange equations in variational problem concerning the functional (3.1) are given by*

$$(3.5) \quad \begin{cases} \frac{d^2 x^i}{dt^2} + 2(G^i(x, y) + H^i(x, y)) = 0, \\ y^i = \frac{dx^i}{dt}, \end{cases}$$

where  $G^i(x, y) = \frac{1}{2} \gamma_{rs}^i y^r y^s$ ,  $H^i(x, y) = \frac{1}{2} \left( \frac{a + 2b\beta}{2} F_h^i y^h + A^i \bar{B} \right)$ , if  $b > 0$ ,  $\bar{B} = \frac{\sqrt{b}}{2} \tilde{F} - \frac{1}{\sqrt{b}(1+c^2)} A_k \gamma_{rs}^k \frac{dx^r}{dt} \frac{dx^s}{dt} - \frac{a + 2b\beta}{2(1+c^2)} A^i F_{ih} \frac{dx^h}{dt} - \frac{b}{2(1+c^2)} \tilde{F} A$ ,  $\tilde{F} = \left( \frac{\partial A_r}{\partial x^s} + \frac{\partial A_s}{\partial x^r} \right) \frac{dx^s}{dt} \frac{dx^r}{dt}$ ,  $A = A^i A_i$  and  $H^i(x, y) = \frac{1}{2} \left( \frac{a + 2b\beta}{2} F_h^i y^h + A^i \underline{B} \right)$ , if  $b < 0$ ,  $d^2 \neq 1$ ,  $\underline{B} = \frac{\sqrt{-b}}{2} \tilde{F} + \frac{1}{\sqrt{-b}(1-d^2)} A_k \gamma_{rs}^k \frac{dx^r}{dt} \frac{dx^s}{dt} + \frac{a + 2b\beta}{2(1-d^2)} A^i F_{ih} \frac{dx^h}{dt} + \frac{b}{2(1-d^2)} \tilde{F} A$ .

The equations (3.5) determine a spray defined only by the Lagrangian  $\mathcal{L}$  from (1.1), so we can develop the geometry of the Lagrange space  $L^n = (M, L)$  using this canonical spray only. We have then:

**Theorem 3.2** *The canonical nonlinear connection of the Lagrange space  $L^n$  is given by:*

$$(3.6) \quad \begin{cases} N_j^i = \underset{(1)}{\overset{o}{N}}_j^i - \underset{(1)}{\overline{A}}_j^i, & \text{if } b > 0, \\ N_j^i = \underset{(1)}{\overset{o}{N}}_j^i - \underset{(1)}{\underline{A}}_j^i, & \text{if } b < 0, \end{cases}$$

where  $\bar{A}_j^i = \frac{1}{2} \frac{\partial}{\partial y^j} \left\{ \frac{a+2b\beta}{2} F_h^i y^h + A^i \bar{B} \right\}$ ,  $\underline{A}_j^i = \frac{1}{2} \frac{\partial}{\partial y^j} \left\{ \frac{a+2b\beta}{2} F_h^i y^h + A^i \underline{B} \right\}$ .

We can prove now:

**Theorem 3.3** The linear N-connection on the Lagrange space  $L^n$  is:

$$(3.7) \quad \left\{ N_j^i = \bar{N}_j^i - \bar{A}_j^i, L_{jk}^i = \bar{F}_{jk}^i + \bar{C}_{jm}^i \bar{A}_k^m + \bar{A}_{jk}^i, C_{jk}^i = \bar{C}_{jk}^i + \bar{C}_{jk}^i \right. \\ \left. (1) \right.$$

in the case  $b > 0$ ,

$$(3.8) \quad \left\{ N_j^i = \bar{N}_j^i - \underline{A}_j^i, L_{jk}^i = \bar{F}_{jk}^i + \bar{C}_{jm}^i \underline{A}_k^m + \underline{A}_{jk}^i, C_{jk}^i = \bar{C}_{jk}^i + \underline{C}_{jk}^i \right. \\ \left. (1) \right.$$

in the case  $b < 0$ , where  $\bar{A}_{jk}^i = \frac{b}{2} g^{ih} \left( \frac{\partial(A_j A_h)}{\partial x^k} + \frac{\partial(A_k A_h)}{\partial x^j} - \frac{\partial(A_j A_k)}{\partial x^h} \right) - \frac{A^i A^h}{2(1+c^2)} \left( \frac{\delta \gamma_{jh}}{\delta x^k} + \frac{\delta \gamma_{kh}}{\delta x^j} - \frac{\delta \gamma_{jk}}{\delta x^h} \right)$ ,  $\underline{A}_{jk}^i = \frac{b}{2} g^{ih} \left( \frac{\partial(A_j A_h)}{\partial x^k} + \frac{\partial(A_k A_h)}{\partial x^j} - \frac{\partial(A_j A_k)}{\partial x^h} \right) + \frac{A^i A^h}{2(1-d^2)} \left( \frac{\delta \gamma_{jh}}{\delta x^k} + \frac{\delta \gamma_{kh}}{\delta x^j} - \frac{\delta \gamma_{jk}}{\delta x^h} \right)$ ,  $\bar{C}_{jk}^i = -\frac{1}{2(1+c^2)} A^i A^h \left( \frac{\partial \gamma_{hj}}{\partial y^k} + \frac{\partial \gamma_{kh}}{\partial y^j} - \frac{\partial \gamma_{jk}}{\partial y^h} \right)$ ,  $\underline{C}_{jk}^i = \frac{1}{2(1-d^2)} A^i A^h \left( \frac{\partial \gamma_{hj}}{\partial y^k} + \frac{\partial \gamma_{kh}}{\partial y^j} - \frac{\partial \gamma_{jk}}{\partial y^h} \right)$ .

Then the whole geometry of the Lagrange space  $L^n$  can be developed only on the base of canonical linear connection given by Theorem 3.3. This connection is a canonical one because it is determined only by the fundamental function (1.1) of the Lagrange space  $L^n = (M, L)$ .

## References

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