

Proper affine vector fields in Bianchi type I Space-Time

Ghulam Shabbir

Abstract. A study of Bianchi type I space-times according to its proper affine vector field is given by using holonomy and decomposability, the rank of the 6×6 Riemann matrix and direct integration techniques. It is shown that the special class of the above space-times admits proper affine vector fields.

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§ 1. Introduction

The aim of this paper is to find the existence of proper affine vector fields in Bianchi type I static space-times by using holonomy and decomposability, the rank of the 6×6 Riemann matrix and direct integration techniques. Through out M is representing the four dimensional, connected, hausdorff space-time manifold with Lorentz metric g of signature $(-, +, +, +)$ The curvature tensor associated with g through Levi-Civita connection Γ , is denoted in component form by $R^a{}_{bcd}$. The usual covariant, partial and Lie derivatives are denoted by a semicolon, a comma and the symbol L , respectively. Round and square brackets denote the usual symmetrization and skew-symmetrization, respectively. The space-time M will be assumed nonflat in the sense that the Riemann tensor does not vanish over any nonempty open subset of M .

A vector field X on M is called an affine vector field if it satisfies

$$(1.1) \quad X_{a;bc} = R_{abcd}X^d.$$

or equivalently,

$$\begin{aligned} X_{a,bc} - \Gamma_{ac}^f X_{f,b} - \Gamma_{bc}^f X_{a,f} - \Gamma_{ab}^e X_{e,c} + \Gamma_{ab}^e \Gamma_{ec}^f X_f - (\Gamma_{ab}^e)_{,c} X_e - \Gamma_{ab}^f \Gamma_{cf}^e X_e \\ + \Gamma_{fb}^e \Gamma_{ca}^f X_e + \Gamma_{af}^e \Gamma_{bc}^f X_e = R_{abcd}X^d. \end{aligned}$$

If one decomposes $X_{a;b}$ on M into its symmetric and skew-symmetric parts

$$(1.2) \quad X_{a;b} = \frac{1}{2}h_{ab} + F_{ab} \quad (h_{ab} = h_{ba}, \quad F_{ab} = -F_{ba})$$

then equation (1.1) is equivalent to

$$(1.3) \quad (i) \quad h_{ab;c} = 0 \quad (ii) \quad F_{ab;c} = R_{abcd}X^d \quad (iii) \quad F_{ab;c}X^c = 0.$$

Such a vector field X is called affine if the local diffeomorphism ϕ_t (for appropriate t) associated with X map geodesics into geodesics. If $h_{ab} = 2cg_{ab}$, $c \in R$ then the vector field X is called homothetic (and Killing if $c = 0$.) The vector field X is said to be proper affine if it is not homothetic vector field and also X is said to be proper homothetic vector field if it is not Killing vector field on M [2]. Define the subspace S_p of the tangent space T_pM to M at p as those $k \in T_pM$ satisfying

$$(1.4) \quad R_{abcd}k^d = 0.$$

§ 2. Affine vector fields

In this section we will briefly discuss when the space-times admit proper affine vector fields for further details see [4].

Suppose that M is a simple connected space-time. Then the holonomy group of M is a connected Lie subgroup of the identity component of the Lorentz group and is thus characterized by its subalgebra in the Lorentz algebra. These have been labeled into fifteen types $R_1 - R_{15}$ [9, 3]. It follows from [4] that the only such space-times which could admit proper affine vector fields are those which admit nowhere zero covariantly constant second order symmetric tensor field h_{ab} and it is known that this forces the holonomy type to be either $R_2, R_3, R_4, R_6, R_7, R_8, R_{10}, R_{11}$ or R_{13} . Here, we will only discuss the space-times which has the holonomy type R_2, R_4, R_7, R_{10} or R_{13} .

First consider the case when M has type R_{13} . Then one can always set up local coordinates (t, x^1, x^2, x^3) on an open set $U = U_1 \times U_2$, where U_1 is a one dimensional timelike submanifold of U coordinatized by t and U_2 is a three dimensional spacelike submanifold of U coordinatized by x^1, x^2, x^3 and where the above product is a metric product and the metric on U is given by [2]

$$(2.5) \quad ds^2 = -dt^2 + g_{\alpha\beta}dx^\alpha dx^\beta, \quad (\alpha, \beta = 1, 2, 3)$$

where $g_{\alpha\beta}$ depends on x^γ , ($\gamma = 1, 2, 3$). The above space-time is clearly 1 + 3 decomposable. The curvature rank of the above space-time is atmost three and there exists a unique nowhere zero vector field $t_a = t_{,a}$ satisfying $t_{a;b} = 0$ and also $t^a t_a = -1$. From the Ricci Identity $R^a_{bcd}t^d = 0$. It follows from [4] that affine vector fields in this case are

$$(2.6) \quad X = (c_1 t + c_2) \frac{\partial}{\partial t} + Y$$

where $c_1, c_2 \in R$ and Y is a homothetic vector field in the induced geometry on each of the three dimensional submanifolds of constant t .

Now consider the situation when M has type R_{10} . The situation is similar to that of previous R_{13} case except that now we have local decomposition is $U = U_1 \times U_2$, where U_1 is a one dimensional spacelike submanifold of U and U_2 is a three dimensional timelike submanifold of U . The space-time metric on U is given by [2]

$$(2.7) \quad ds^2 = dx^2 + g_{\alpha\beta}dx^\alpha dx^\beta, \quad (\alpha, \beta = 0, 2, 3)$$

where $g_{\alpha\beta}$ depends on x^γ , ($\gamma = 0, 2, 3$). The above space-time is clearly 1 + 3 decomposable. The curvature rank of the above space-time is atmost three and there exists a unique nowhere zero vector field $x_a = x_{,a}$ satisfying $x_{a;b} = 0$ and also $x^a x_a = 1$. From the Ricci Identity $R^a{}_{bcd} x^d = 0$. It follows from [4] that affine vector fields in this case are

$$(2.8) \quad X = (c_1 x + c_2) \frac{\partial}{\partial x} + Y$$

where $c_1, c_2 \in R$ and Y is a homothetic vector field in the induced geometry on each of the three dimensional submanifolds of constant x .

Next suppose M has type R_7 . Then each $p \in M$ has a neighborhood U which decomposes metrically as $U = U_1 \times U_2$, where U_1 is a two dimensional submanifold of U with an induced metric of Lorentz signature and U_2 is a two dimensional submanifold of U with positive definite induced metric. The space-time metric on U is given by [2]

$$(2.9) \quad ds^2 = P_{AB} dx^A dx^B + Q_{\alpha\beta} dx^\alpha dx^\beta,$$

where $P_{AB} = P_{AB}(x^C)$, for all $A, B, C = 0, 1$ and $Q_{\alpha\beta} = Q_{\alpha\beta}(x^\gamma)$, for all $\alpha, \beta, \gamma = 2, 3$ and the above space-time is clearly 2 + 2 decomposable. The space-time (2.8) admits two recurrent vector fields [1] l and n i.e. $l_{a;b} = l_a p_b$ and $n_{a;b} = n_a p_b$, where p_b is the recurrent 1-form. It also admits two covariantly constant second order symmetric tensors which are $2l_{(a} n_{b)}$ and $(x_a x_b + y_a y_b)$. The rank of the 6×6 Riemann matrix is two. It follows from [4] that if X is an affine vector field on M then X decomposes as

$$(2.10) \quad X = X_1 + X_2,$$

where the vector fields X_1 and X_2 are tangent to the two dimensional timelike and spacelike submanifolds, respectively. It also follows from [4] that X_1 and X_2 are homothetic vector fields in their respective submanifolds with their induced geometry. Conversely, every pair of affine vector fields, one in the timelike submanifolds and one spacelike submanifolds give rise to a affine vector field in space-time.

Now suppose that M has type R_4 . Then each $p \in M$ has a neighborhood U which decomposes metrically as $U = U_1 \times U_2 \times U_3$, where U_1 and U_2 are one dimensional submanifold of U and U_3 is a two dimensional submanifold of U . The space-time metric on U is given by [4]

$$(2.11) \quad ds^2 = -dt^2 + dx^2 + g_{AB} dx^A dx^B, \quad (A, B = 2, 3)$$

where g_{AB} depends only on x^C ($C = 2, 3$). The above space-time is clearly 1 + 1 + 2 decomposable. The curvature rank of the above space-time is one and there exist two independent nowhere zero unit timelike and spacelike covariantly constant vector field $t_a = t_{,a}$ and $x_a = x_{,a}$ satisfying $t_{a;b} = 0$ and $x_{a;b} = 0$. From the Ricci identity $R^a{}_{bcd} t_a = 0$ and $R^a{}_{bcd} x_a = 0$. It follows from [4] that affine vector fields in this case are

$$(2.12) \quad X = (c_1 t + c_2 x + c_3) \frac{\partial}{\partial t} + (c_4 t + c_5 x + c_6) \frac{\partial}{\partial x} + Y$$

where $c_1, c_2, c_3, c_4, c_5, c_6 \in R$ and Y is a homothetic vector field in the induced geometry on each of the two dimensional submanifolds of constant t and x .

Now suppose that M has type R_2 . Here each $p \in M$ admits a neighborhood U which decomposes metrically as $U = U_1 \times U_2 \times U_3$, where U_1 and U_2 are one dimensional submanifold of U and U_3 is a two dimensional submanifold of U . The space-time metric on U is given by [4]

$$(2.13) \quad ds^2 = dy^2 + dz^2 + g_{AB}dx^A dx^B, \quad (A, B = 0, 1)$$

where g_{AB} depends only on x^C ($C = 0, 1$). The above space-time is clearly $1 + 1 + 2$ decomposable. The curvature rank of the above space-time is one and there exist two independent nowhere zero unit spacelike covariantly constant vector fields $y_a = y_{,a}$ and $z_a = z_{,a}$ satisfying $y_{a;b} = 0$ and $z_{a;b} = 0$. From the Ricci identity $R^a_{bcd}y_a = 0$ and $R^a_{bcd}z_a = 0$. It follows from [4] that affine vector fields in this case are

$$(2.14) \quad X = (c_1y + c_2z + c_3)\frac{\partial}{\partial y} + (c_4y + c_5z + c_6)\frac{\partial}{\partial z} + Y$$

where $c_1, c_2, c_3, c_4, c_5, c_6 \in R$ and Y is a homothetic vector field in the induced geometry on each of the two dimensional submanifolds of constant y and z .

§ 3. Main results

As mentioned in section 2, the space-times which can admit proper affine vector fields having holonomy type $R_2, R_3, R_4, R_6, R_7, R_8, R_{10}, R_{11}$ or R_{13} . It also follows from [1] that the rank of the 6×6 Riemann matrix is atmost three. Here in this paper we will consider the rank of the 6×6 Riemann matrix to study affine vector fields in spherically symmetric static space-time. Consider a spherically symmetric static space-time in the usual coordinate system (t, x, y, z) with line element [8]

$$(3.15) \quad ds^2 = -dt^2 + f(t)dr^2 + k(t)dy^2 + h(t)dz^2$$

where f, k and h are some nowhere zero functions of t only. It follows from [6], the above space-time admits three independent Killing vector fields which are

$$\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}.$$

The non-zero independent components of the Riemann tensor are [6]

$$(3.16) \quad \begin{aligned} R_{1010} &= -\frac{1}{4} \frac{2f\ddot{f} - \dot{f}^2}{f} = \alpha_1 & R_{2020} &= -\frac{1}{4} \frac{2k\ddot{k} - \dot{k}^2}{k} = \alpha_2 \\ R_{3030} &= -\frac{1}{4} \frac{2h\ddot{h} - \dot{h}^2}{h} = \alpha_3 & R_{2121} &= \frac{\dot{k}\dot{f}}{4} = \alpha_4 \\ R_{3131} &= \frac{\dot{h}\dot{f}}{4} = \alpha_5 & R_{2121} &= \frac{\dot{k}\dot{h}}{4} = \alpha_6 \end{aligned}$$

Writing the curvature tensor with components R_{abcd} at p as a 6×6 symmetric matrix in a well known way [7]

$$(3.17) \quad R_{abcd} = \text{diag}(\alpha_1, \alpha_2, \alpha_3, \alpha_4, \alpha_5, \alpha_6)$$

where $\alpha_1, \alpha_2, \alpha_3, \alpha_4, \alpha_5$ and α_6 are real functions of t . The 6-dimensional labeling is in the order 01, 02, 03, 12, 13, 23 with $x^0 = t$. We are only interested in those cases when the rank of the 6×6 Riemann matrix is less than or equal to three (excluding the flat cases). We thus obtain the following possibilities:

- (A1) Rank = 3, $f \in R - \{0\}, k = k(t), h = h(t)$
- (A2) Rank = 3, $k \in R - \{0\}, f = f(t), h = h(t)$
- (A3) Rank = 3, $h \in R - \{0\}, k = k(t), f = f(t)$
- (A4) Rank = 3, $f = f(t), k = k(t), h = h(t), 2k\ddot{k} - \dot{k}^2 = 0, 2h\ddot{h} - \dot{h}^2 = 0, 2f\ddot{f} - \dot{f}^2 = 0$
- (B1) Rank = 2, $f \in R - \{0\}, k = k(t), h = h(t), 2k\ddot{k} - \dot{k}^2 = 0$
- (B2) Rank = 2, $k \in R - \{0\}, f = f(t), h = h(t), 2h\ddot{h} - \dot{h}^2 = 0$
- (B3) Rank = 2, $h \in R - \{0\}, f = f(t), k = k(t), 2f\ddot{f} - \dot{f}^2 = 0$
- (B4) Rank = 2, $f \in R - \{0\}, k = k(t), h = h(t), 2h\ddot{h} - \dot{h}^2 = 0$
- (B5) Rank = 2, $k \in R - \{0\}, f = f(t), h = h(t), 2f\ddot{f} - \dot{f}^2 = 0$
- (B6) Rank = 2, $k \in R - \{0\}, f = f(t), k = k(t), 2k\ddot{k} - \dot{k}^2 = 0$
- (C1) Rank = 1, $f, h \in R - \{0\}, k = k(t)$
- (C2) Rank = 1, $k, h \in R - \{0\}, f = f(t)$
- (C3) Rank = 1, $f, k \in R - \{0\}, h = h(t)$
- (D1) Rank = 1, $f \in R - \{0\}, k = k(t), h = h(t), 2k\ddot{k} - \dot{k}^2 = 0, 2h\ddot{h} - \dot{h}^2 = 0$
- (D2) Rank = 1, $k \in R - \{0\}, f = f(t), h = h(t), 2f\ddot{f} - \dot{f}^2 = 0, 2h\ddot{h} - \dot{h}^2 = 0$
- (D3) Rank = 1, $h \in R - \{0\}, k = k(t), f = f(t), 2k\ddot{k} - \dot{k}^2 = 0, 2f\ddot{f} - \dot{f}^2 = 0$

We will consider each case in turn.

Case A1

In this case $f \in R - \{0\}, k = k(t), h = h(t)$ and the rank of the 6×6 Riemann matrix is 3 and there exists a unique (up to a multiple) nowhere zero spacelike vector field $x_a = x_{,a}$ satisfying $x_{a;b} = 0$ (and so, from the Ricci identity $R^a{}_{bcd}x_a = 0$). The line element can, after a recaling of x , be written in the form

$$(3.18) \quad ds^2 = dx^2 + (-dt^2 + kdy^2 + hdz^2).$$

The space-time is clearly 1+3 decomposable and its holonomy type is R_{10} . The affine vector fields in this case [4] are

$$(3.19) \quad X = (c_3x + c_4)\frac{\partial}{\partial x} + X'$$

where $c_3, c_4 \in R$ and X' is a homothetic vector field in the induced geometry on each of the three dimensional submanifolds of constant x . The completion of case A necessities finding an homothetic vector fields in the induced geometry of the submanifolds of constant x . The induced metric $g_{\alpha\beta}$ (where $\alpha, \beta = 0, 2, 3$) with nonzero components is given by

$$(3.20) \quad g_{00} = -1, \quad g_{22} = k(t), \quad g_{33} = h(t).$$

A vector field X' is a homothetic vector field if it satisfies $L_{X'}g_{\alpha\beta} = 2cg_{\alpha\beta}$, where $c \in R$. One can expand by using (3.20) to get

$$(3.21) \quad X^0{}_{,0} = c$$

$$(3.22) \quad kX^2_{,0} - X^0_{,2} = 0$$

$$(3.23) \quad hX^3_{,0} - X^0_{,3} = 0$$

$$(3.24) \quad \dot{k}X^0 + 2kX^2_{,2} = 2kc$$

$$(3.25) \quad kX^2_{,3} + hX^3_{,2} = 0$$

$$(3.26) \quad \dot{h}X^0 + 2hX^3_{,3} = 2hc.$$

Equation (3.21), (3.22) and (3.23) give

$$X^0 = ct + A^1(y, z), \quad X^2 = A^1_y(y, z) \int \frac{dt}{k} + A^2(y, z),$$

$$X^3 = A^1_z(y, z) \int \frac{dt}{h} + A^3(y, z),$$

where $A^1(y, z)$, $A^2(y, z)$ and $A^3(y, z)$ are functions of integration. If one proceeds further after a strightforward calculation one can find that proper homothetic vector fields exist if and only if

$$(3.27) \quad k = (at + c)^2, \quad h = (bt + d)^2$$

where $a, b, c, d \in R(a, b \neq 0)$. Substituting (3.27) in (3.16) one finds that the rank of the 6×6 Riemann matrix is reduces to one thus giving a contradiction (since we are assuming that that the rank of the 6×6 Riemann matrix is three). So the only homothetic vector fields in the induced gemetry are the Killing vector fields which are

$$(3.28) \quad X^0 = 0, \quad X^2 = c_1, \quad X^3 = c_2$$

where $c_1, c_2 \in R$. The affine vector fields in this case are (from (3.19) and (3.28))

$$(3.29) \quad X^0 = 0, \quad X^1 = xc_3 + c_4, \quad X^2 = c_1, \quad X^3 = c_2$$

The above space-time (3.18) admits four independent affine vector fields in which three are Killing vector fields and one is proper affine vector field. Cases (A2) and (A3) are exactly same.

Case A4

In this case $f = f(t)$, $k = k(t)$, $h = h(t)$, $2k\ddot{k} - \dot{k}^2 = 0$, $2h\ddot{h} - \dot{h}^2 = 0$ and $2f\ddot{f} - \dot{f}^2 = 0$. Equations $2k\ddot{k} - \dot{k}^2 = 0$, $2h\ddot{h} - \dot{h}^2 = 0$ and $2f\ddot{f} - \dot{f}^2 = 0$ imply $k = (a_1t + a_2)^2$, $h = (a_3t + a_4)^2$ and $f = (a_5t + a_6)^2$, respectively, where $a_1, a_2, a_3, a_4, a_5, a_6 \in R(a_1, a_3, a_5 \neq 0)$. We first suppose that $a_1 \neq a_3$, $a_1 \neq a_5$, $a_3 \neq a_5$, $a_2 \neq a_4$, $a_2 \neq a_6$ and $a_4 \neq a_6$. The rank of the 6×6 Riemann matrix is 3 and there exists a unique (up to a multiple) solution $t_a = t_{,a}$ of equation (1.4) satisfying but t_a is not covariantly constant. The line element is

$$(3.30) \quad ds^2 = -dt^2 + (a_5t + a_6)^2 dx^2 + (a_1t + a_2)^2 dy^2 + (a_3t + a_4)^2 dz^2.$$

Affine vector fields in this case are

$$(3.31) \quad X^0 = 0, \quad X^1 = c_3, \quad X^2 = c_1, \quad X^3 = c_2,$$

where $c_1, c_2, c_3 \in R$. Affine vector fields in this case are Killing vector fields. Now $a_1 = a_3 = a_5 = a(say) \in R(a \neq 0)$ and $a_2 = a_4 = a_6 = b \in R$. The line element is

$$(3.32) \quad ds^2 = -dt^2 + (at + b)^2(dx^2 + dy^2 + dz^2).$$

It follows from [10] that affine vector fields in this case are

$$X^0 = tc_7, \quad X^1 = -yc_2 + zc_3 + c_4,$$

$$(3.33) \quad X^2 = xc_2 - yc_5 + c_6, \quad X^3 = -xc_3 + yc_5 + c_1,$$

where $c_1, c_2, c_3, c_4, c_5, c_6, c_7 \in R$. The above space-time (3.32) admits seven independent affine vector fields in which six are Killing vector fields and one is proper affine vector field.

Case B1

In this case $f \in R - \{0\}, k = k(t), h = h(t)$ and $2k\ddot{k} - \dot{k}^2 = 0$. Equation $2k\ddot{k} - \dot{k}^2 = 0$ imply $k = (at + b)^2$, where $a, b \in R(a \neq 0)$. The rank of the 6×6 Riemann matrix is two and there exists a unique (up to a multiple) nowhere zero spacelike vector field $x_a = x_{,a}$ satisfying $x_{a;b} = 0$ (and so, from the Ricci identity $R^a{}_{bcd}x_a = 0$). After a suitable recaling of x , the line element takes the form

$$(3.34) \quad ds^2 = dx^2 + (-dt^2 + (at + b)^2dy^2 + h dz^2).$$

The space-time is clearly 1+3 decomposable and its holonomy type is R_{10} . The affine vector fields in this case are of the form (3.19). The completion of case B1 necessities finding an homothetic vector fields in the induced geometry of the submanifolds of constant x . The induced metric $g_{\alpha\beta}$ (where $\alpha, \beta = 0, 2, 3$) with nonzero components is given by

$$(3.35) \quad g_{00} = -1, \quad g_{22} = (at + b)^2, \quad g_{33} = h(t).$$

If one proceed further exactly on the same lines as we did in case A1, one can find that the homothetic vector fields in the induced geometry are the Killing vector fields which are given in (3.28). Affine vector fields in this case are given in (3.29). Cases (B2), (B3), (B4), (B5) and (B6) are exactly same.

Case C1

In this case $f, h \in R - \{0\}, k = k(t)$ and the rank of the 6×6 Riemann matrix is one. There exist two independent nowhere zero spacelike vector fields $x_a = x_{,a}$ and $z_a = z_{,a}$ satisfying $x_{a;b} = 0$ and $z_{a;b} = 0$. From the Ricci identity $R^a{}_{bcd}x_a = R^a{}_{bcd}z_a = 0$. The line element can, after a suitable recaling of x and z , be written as

$$(3.36) \quad ds^2 = dx^2 + dz^2 + (-dt^2 + k dy^2).$$

The space-time is clearly 1 + 1 + 2 decomposable and its holonomy type is R_{10} . Affine vector fields in this case are [4]

$$(3.37) \quad X = (c_1x + c_2z + c_3) \frac{\partial}{\partial x} + (c_4x + c_5z + c_6) \frac{\partial}{\partial z} + X'$$

where $c_1, c_2, c_3, c_4, c_5, c_6 \in R$ and X' is a homothetic vector field in the induced geometry on each of the two dimensional submanifolds of constant x and z . The next step is to find the homothetic vector fields in the induced geometry of the submanifolds of constant x and z . The induced metric g_{AB} (where $A, B = 0, 2$) with nonzero components are given by

$$(3.38) \quad g_{00} = -1, \quad g_{22} = k(t).$$

A vector field X' is a homothetic vector field if it satisfies $L_{X'} g_{AB} = 2c g_{AB}$, where $c \in R$. One can expand by using (3.38) to get

$$(3.39) \quad X^0_{,0} = c$$

$$(3.40) \quad kX^2_{,0} - X^0_{,2} = 0$$

$$(3.41) \quad \dot{k}X^0 - +2kX^2_{,2} = 2ck$$

Equation (3.39) gives $X^0 = ct + A^1(y)$, where $A^1(y)$ is a function of integration. Using X^0 in equation (3.40) we get $X^2 = A^1_y(y) \int \frac{dt}{k} + A^2(y)$, where $A^2(y)$ is a function of integration. If one proceeds further, after a straightforward calculation one finds that proper homothetic vector field exist if and only if $k = at^2$, where $a \in R - \{0\}$. Substituting the value of k into (3.16), one finds that the rank of the 6×6 Riemann matrix reduces to zero thus giving a contradiction (since we are assuming that the rank of the 6×6 Riemann matrix is one). So homothetic vector fields in the induced geometry of constant x and z are Killing vector fields. If one proceeds further one finds there exist two possibilities:

$$(a) k\left(\frac{\dot{k}}{2k}\right) = n \quad (b) k\left(\frac{\dot{k}}{2k}\right) \neq n$$

where $n \in R$.

Case C1a

In this case further three possibilities exist

$$(i) n > 0, \quad (ii) n < 0, \quad (iii) n = 0$$

We will consider each case in turn.

(i) Affine vector fields in this case are

$$(3.42) \quad \begin{aligned} X^0 &= c_7 \sin \sqrt{n}y + c_8 \cos \sqrt{n}y, \\ X^1 &= c_1y + c_2z + c_3, \\ X^2 &= \sqrt{n}(c_7 \cos \sqrt{n}y - c_8 \sin \sqrt{n}y) \int \frac{dt}{k} + c_9, \\ X^3 &= c_4y + c_5z + c_6, \end{aligned}$$

provided that $k(\frac{\dot{k}}{2k}) = n$, where $c_7, c_8, c_9 \in R$.

(ii) In this case $n < 0$. Put $n = -N$, where $N \in R(N > 0)$. Affine vector fields in this case are

$$(3.43) \quad \begin{aligned} X^0 &= c_7 \sinh \sqrt{N}y + c_8 \cosh \sqrt{N}y, \\ X^1 &= c_1y + c_2z + c_3, \\ X^2 &= \sqrt{N}(c_7 \cosh \sqrt{N}y + c_8 \sinh \sqrt{N}y) \int \frac{dt}{k} + c_9, \\ X^3 &= c_4y + c_5z + c_6, \end{aligned}$$

provided that $k(\frac{\dot{k}}{2k}) = -N$, where $c_7, c_8, c_9 \in R$.

(iii) In this case $n = 0$ implies $k = e^{at+b}$, where $a, b \in R(a \neq 0)$. Affine vector fields in this case are

$$(3.44) \quad \begin{aligned} X^0 &= c_7y + c_8, \\ X^1 &= c_1y + c_2z + c_3, \\ X^2 &= -\frac{c_7}{a}e^{-(at+b)} + c_9, \\ X^3 &= c_4y + c_5z + c_6, \end{aligned}$$

where $c_7, c_8, c_9 \in R$.

(C1b) Affine vector fields in this case are

$$(3.45) \quad X^0 = c_7, X^1 = c_1y + c_2z + c_3, X^2 = 0, X^3 = c_4y + c_5z + c_6,$$

where $c_7 \in R$. This completes case C1. Cases C2 and C3 are exactly same.

(D1)

In this case $f \in R - \{0\}$, $k = k(t)$, $h = h(t)$, $2k\ddot{k} - \dot{k}^2 = 0$ and $2h\ddot{h} - \dot{h}^2 = 0$. Equations $2k\ddot{k} - \dot{k}^2 = 0$ and $2h\ddot{h} - \dot{h}^2 = 0$ imply $k = (at + b)^2$ and $h = (ct + d)^2$ respectively, where $a, b, c, d \in R(a \neq c, a, c \neq 0)$. The rank of the 6×6 Riemann matrix is one and there exist two independent nowhere zero solutions $t_a = t_{,a}$ and $x_a = x_{,a}$ of equation (1.4) with t timelike and x spacelike vector fields, respectively and $x_{a,b} = 0$. After a rescaling of x the line element is

$$(3.46) \quad ds^2 = -dt^2 + dx^2 + (at + b)^2 dy^2 + (ct + d)^2 dz^2.$$

The space-time is clearly 1 + 3 decomposable and the rank of the 6×6 Riemann matrix is one. Substituting the above information into affine equations and after a straightforward calculation one find affine vector fields in this case are

$$(3.47) \quad X^0 = 0, X^1 = c_3t + c_4x + c_5, X^2 = c_1, X^3 = c_2,$$

where $c_1, c_2, c_3, c_4, c_5 \in R$.

Now suppose that $a = c, b = d \in R(a, c \neq 0)$. The line element is

$$(3.48) \quad ds^2 = -dt^2 + dx^2 + (at + b)^2(dy^2 + dz^2).$$

We know from [7, 5] that affine vector fields in this case are

$$(3.49) \quad X^0 = c_4(at + b), X^1 = c_5x + c_6, X^2 = -c_3z + c_1, X^3 = c_3y + c_2,$$

where $c_1, c_2, c_3, c_4, c_5, c_6 \in R$. Cases (D2) and (D3) are exactly the same.

Summary

In this paper a study of Bianchi type I space-times according to their proper affine vector fields is given. An approach is developed to study proper affine vector fields in the above space-times by using the rank of the 6×6 Riemann matrix and holonomy. From the above study we obtain the following results:

- (i) The case when the rank of the 6×6 Riemann matrix is three and there exists a nowhere zero independent spacelike vector field which is the solution of equation (1.4) and is not covariantly constant. This is the space-time (3.30) and it admits affine vector fields which are Killing vector fields (for details see Case A4).
- (ii) The case when the rank of the 6×6 Riemann matrix is one and there exist two nowhere zero independent solution of equation (1.4) but only one independent covariantly constant vector field. This is the space-time (3.48) and it admits proper affine vector fields (see Case D1).
- (iii) The case when the rank of the 6×6 Riemann matrix is two or three and there exists a nowhere zero independent spacelike vector field which is the solution of equation (1.4) and also covariantly constant. This is the space-time (3.18) and (3.34) and it admits proper affine vector fields (see Cases A1 and B1).
- (iv) In the case when the rank of the 6×6 Riemann matrix one there exist two nowhere zero independent spacelike and timelike vector fields which are solutions of equation (1.4) and are covariantly constant. This is the space-time (3.36) and it admits proper affine vector fields (see Case C1).
- (v) The case when the rank of the 6×6 Riemann matrix is three and there exists a nowhere zero independent spacelike vector field which is the solution of equation (1.4) and is not covariantly constant. This is the space-time (3.32) and it admits proper affine vector fields (see equation (3.33)).

References

- [1] G. S. Hall and W. Kay, *Curvature structure in general relativity*, Journal of Mathematical Physics, 29 (1988), 420-427; G. S. Hall and W. Kay, *Holonomy groups in general relativity*, Journal of Mathematical Physics, 29 (1988), 428-432.
- [2] G. S. Hall and J. da. Costa, *Affine collineations in space-time*, Journal of Mathematical Physics, 29 (1988), 2465-2472.
- [3] G. S. Hall, *Covariantly constant tensors and holonomy structure in general relativity*, Journal of Mathematical Physics, 32 (1991), 181-187.
- [4] G. S. Hall, D. J. Low and J. R. Pulham, *Affine collineations in general relativity and their fixed point structure*, Journal of Mathematical Physics, 35 (1994), 5930-5944.
- [5] G. S. Hall and D. P. Lonie, *projective collineations in space-times*, Classical and Quantum Gravity, 12 (1995), 1007-1020.
- [6] G. Shabbir, *Proper curvature collineations in Bianchi type I space-times*, Gravitation and Cosmology, 9 (2003), 139-145.
- [7] G. Shabbir, *Proper projective symmetry in plane symmetric static space-times*, Classical and Quantum Gravity, 21 (2004), 339-347.

- [8] H. Stephani, D. Kramer, M. A. H. MacCallum, C. Hoenselears and E. Herlt, *exact solutions of Einstein's field equations*, Cambridge University Press, 2003.
- [9] J. F. Schell, *Classification of four-dimensional Riemannian spaces*, Journal of Mathematical Physics, 2 (1961), 202-205.
- [10] R. Maartens, *Affine collineations in Robertson-Walker space-time*, Journal of Mathematical Physics, 29 (1987), 2051-2052.

Author's address:

Ghulam Shabbir
Faculty of Engineering Sciences,
GIK Institute of Engineering Sciences and Technology,
Topi, Swabi, NWFP, Pakistan
e-mail: shabbir@giki.edu.pk